

The force-free magnetosphere of a rotating black hole

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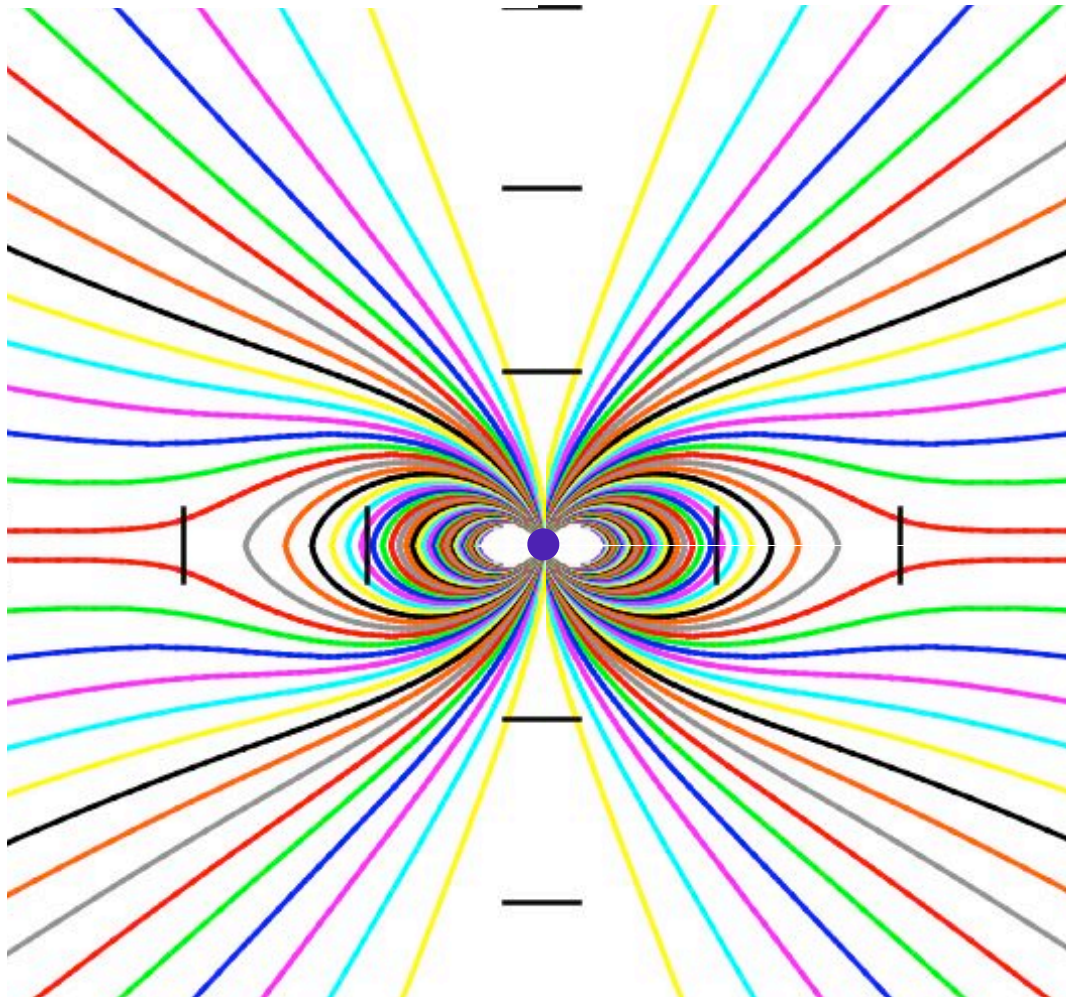
- ...Goldreich & Julian laid the foundations for the theory of 'pulsar electrodynamics' in 1969...
- ...A compelling mechanism by which black holes can energize quasars was not found until 13 years after the quasar discovery... The long delay is surprising, since the mechanism is essentially the same as in the pulsar case: magnetic fields, embedded in a rotating black hole and a surrounding accretion disk, transmit rotational and orbital energy to distant radiating particles...
- Many astrophysicists feel uncomfortable in curved spacetime...

- Macdonald & Thorne reformulated electrodynamics with the hope that it may catalyze pulsar-experienced astrophysicists to begin research on black-hole electrodynamics and to bring to bear on this topic their lore about the 'axisymmetric pulsar problem'...

Topics

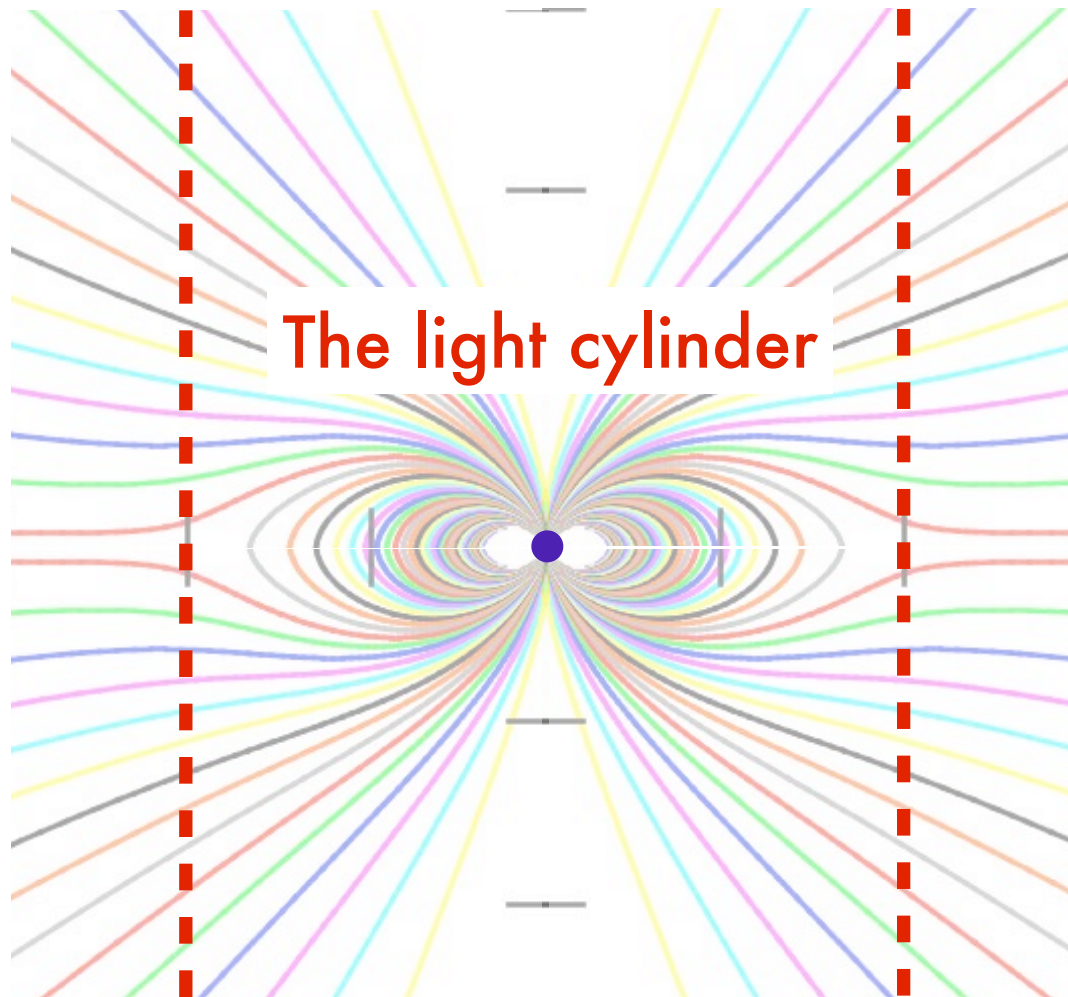
- The analogy with pulsars
- The GR pulsar equation
 - The generalized light cylinders
 - Solution
- Black hole jets
- The Cosmic Battery

The analogy with pulsars



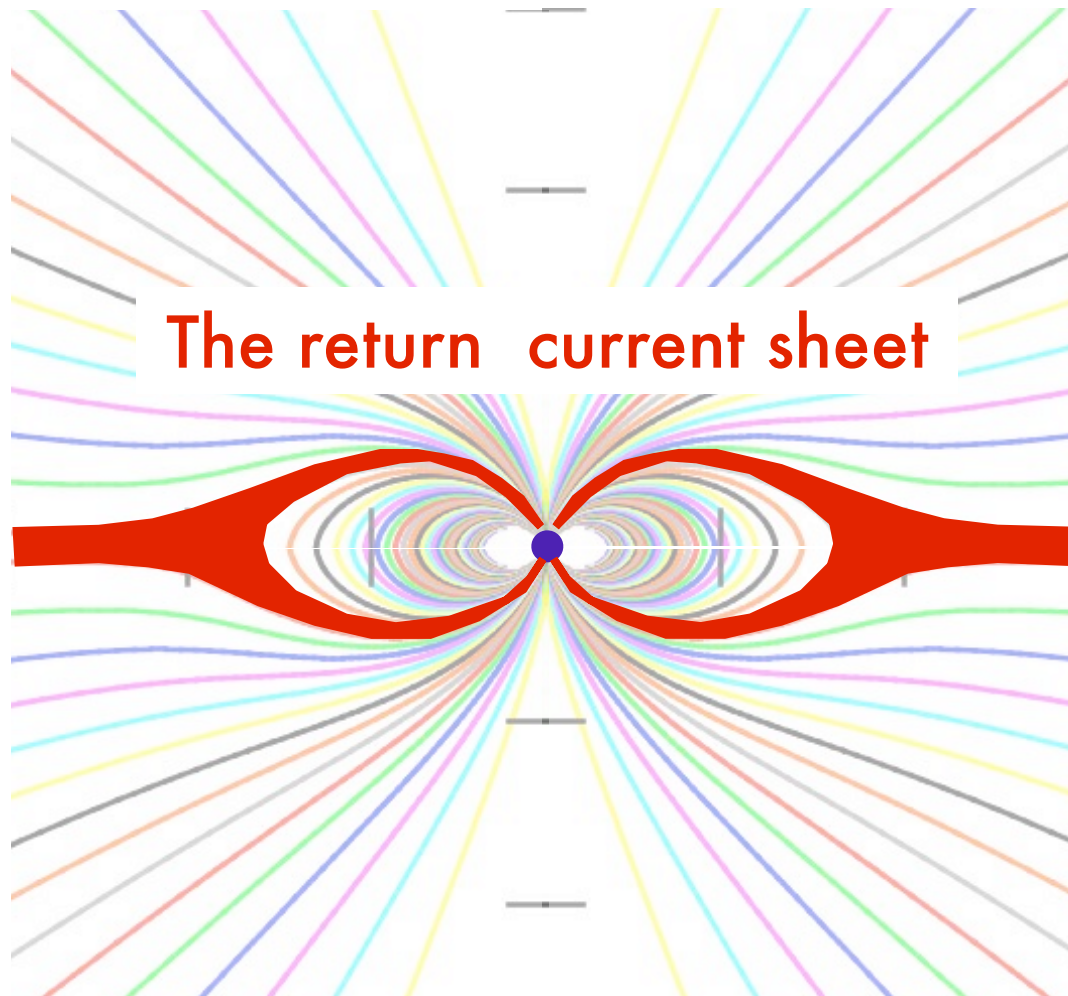
Contopoulos, Kazanas & Fendt 1999

The analogy with pulsars



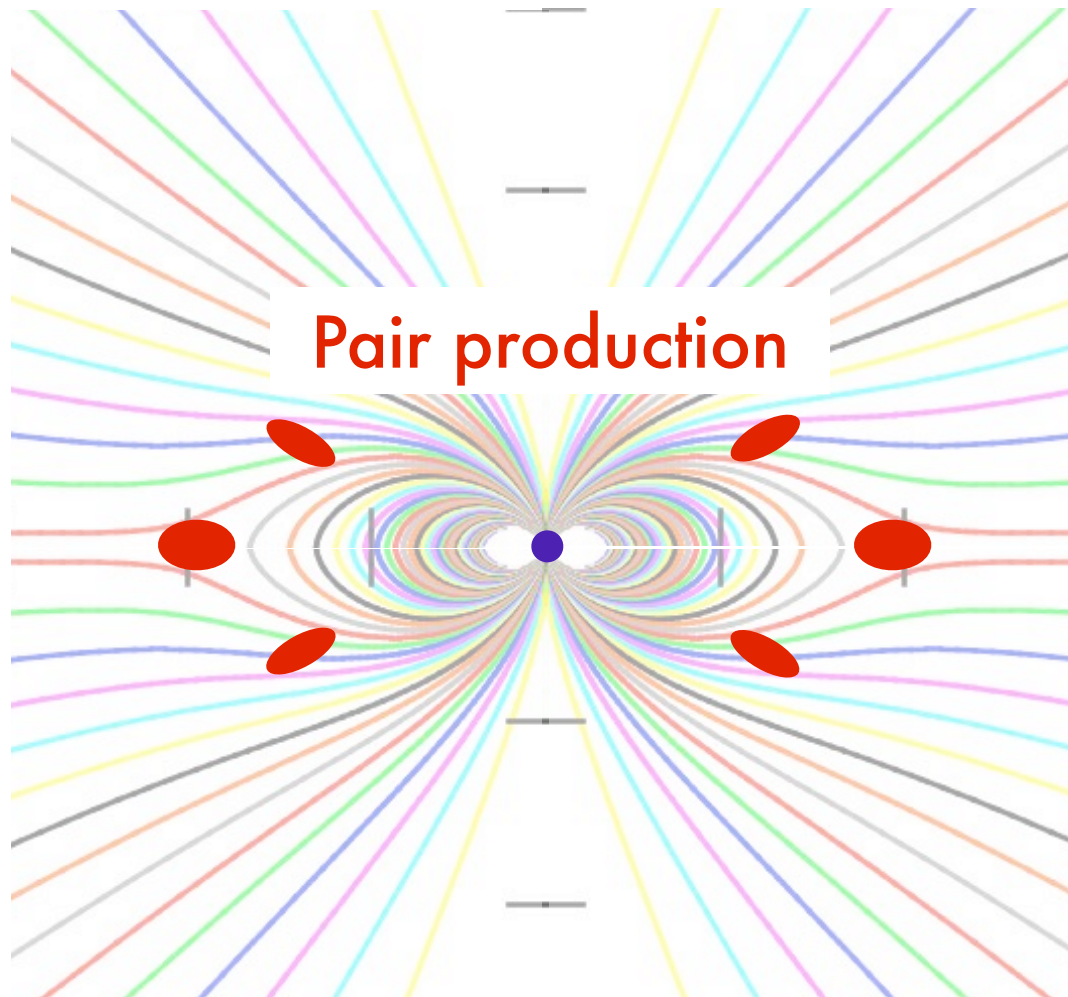
Contopoulos, Kazanas & Fendt 1999

The analogy with pulsars



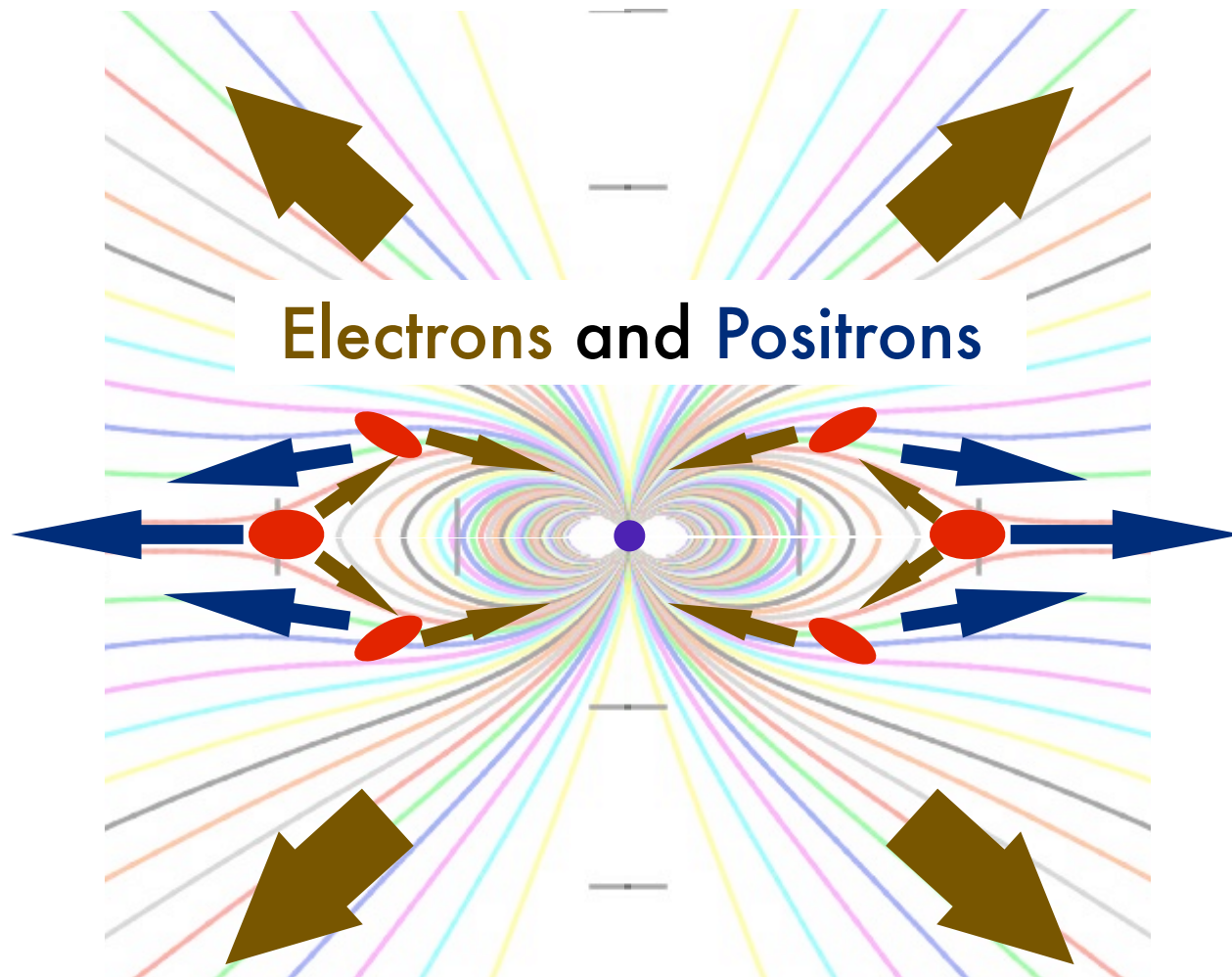
Contopoulos, Kazanas & Fendt 1999

The analogy with pulsars



Contopoulos, Kazanas, Harding, Kalapotharakos

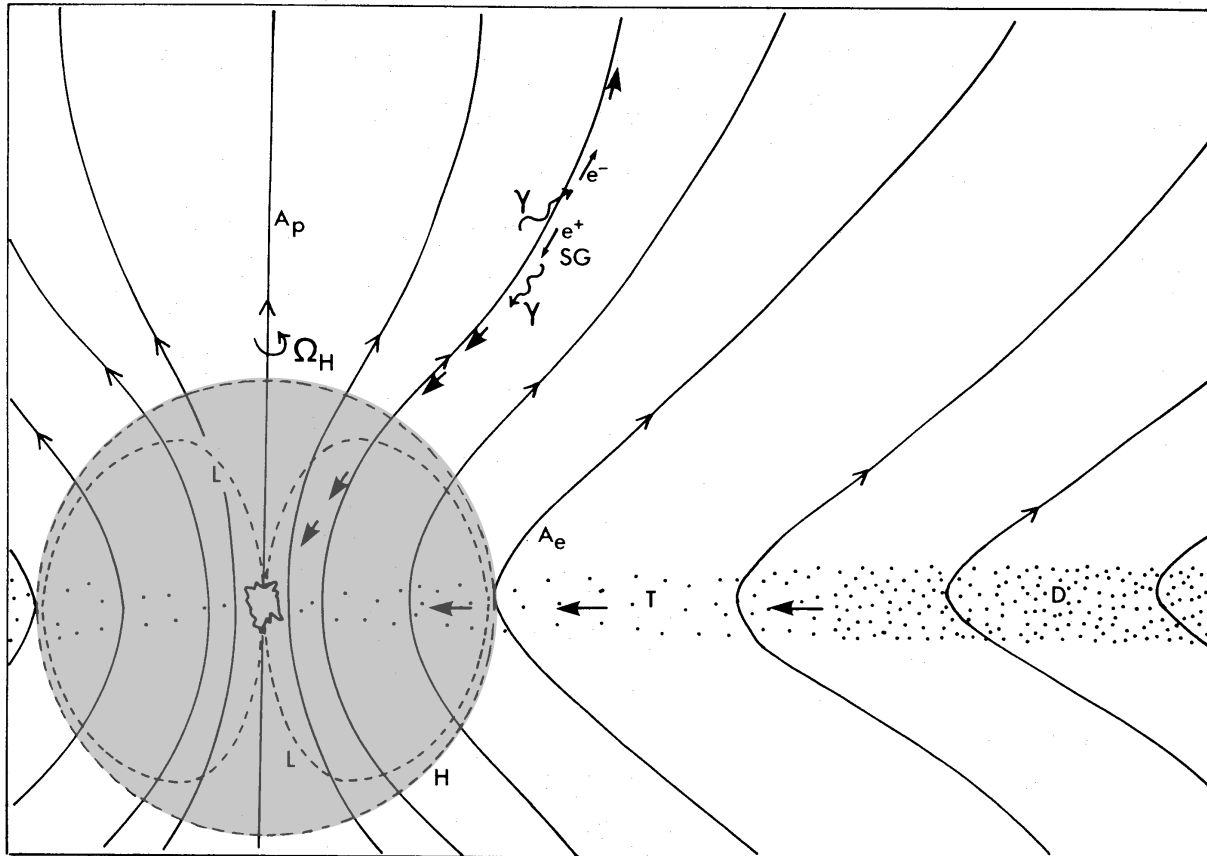
The analogy with pulsars



Contopoulos, Kazanas, Harding, Kalapotharakos

Blandford-Znajek revisited

$$\mathcal{E}_{EM} \propto \omega(\Omega_{\text{BH}} - \omega)\Psi_m^2 \sim \Omega_{\text{BH}}^2 \Psi_m^2$$



Blandford, Znajek 1977

Blandford-Znajek revisited

- Radio loud / radio quiet AGN
- Jet formation and disruption in X-ray binaries
- No relation between BH spin and jet power?!!

The GR pulsar equation

$$\begin{aligned} ds^2 &= g_{\mu\nu} dx^\mu dx^\nu \\ &= -\alpha^2 dt^2 + \frac{A \sin^2 \theta}{\Sigma} (d\phi - \Omega dt)^2 + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 \end{aligned}$$

The GR pulsar equation

$$\begin{aligned} ds^2 &= g_{\mu\nu} dx^\mu dx^\nu \\ &= -\alpha^2 dt^2 + \frac{A \sin^2 \theta}{\Sigma} (d\phi - \Omega dt)^2 + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 \end{aligned}$$

$$\tilde{\nabla} \cdot \tilde{B} = 0$$

$$\tilde{\nabla} \cdot \tilde{E} = 4\pi\rho_e$$

$$\tilde{\nabla} \times (\alpha\tilde{B}) = 4\pi\alpha\tilde{J}$$

$$\nabla \times (\alpha\tilde{E}) = 0.$$

$$\tilde{E} \cdot \tilde{B} = 0$$

The GR pulsar equation

$$\begin{aligned} ds^2 &= g_{\mu\nu} dx^\mu dx^\nu \\ &= -\alpha^2 dt^2 + \frac{A \sin^2 \theta}{\Sigma} (d\phi - \Omega dt)^2 + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 \end{aligned}$$

$$\begin{aligned} \tilde{\nabla} \cdot \tilde{B} &= 0 \\ \tilde{\nabla} \cdot \tilde{E} &= 4\pi\rho_e \\ \tilde{\nabla} \times (\alpha\tilde{B}) &= 4\pi\alpha\tilde{J} \\ \nabla \times (\alpha\tilde{E}) &= 0. \\ \tilde{E} \cdot \tilde{B} &= 0 \end{aligned} \quad \tilde{B}(r, \theta) = \frac{1}{\sqrt{A} \sin \theta} \left\{ \Psi_{,\theta}, -\sqrt{\Delta} \Psi_{,r}, \frac{2I\sqrt{\Sigma}}{\alpha} \right\}$$
$$\tilde{E}(r, \theta) = \frac{\Omega - \omega}{\alpha\sqrt{\Sigma}} \left\{ \sqrt{\Delta} \Psi_{,r}, \Psi_{,\theta}, 0 \right\} .$$

The GR pulsar equation

$$\begin{aligned} ds^2 &= g_{\mu\nu} dx^\mu dx^\nu \\ &= -\alpha^2 dt^2 + \frac{A \sin^2 \theta}{\Sigma} (d\phi - \Omega dt)^2 + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 \end{aligned}$$

$$\begin{aligned} \tilde{\nabla} \cdot \tilde{B} &= 0 \\ \tilde{\nabla} \cdot \tilde{E} &= 4\pi\rho_e \\ \tilde{\nabla} \times (\alpha\tilde{B}) &= 4\pi\alpha\tilde{J} \\ \nabla \times (\alpha\tilde{E}) &= 0. \\ \tilde{E} \cdot \tilde{B} &= 0 \end{aligned} \quad \tilde{B}(r, \theta) = \frac{1}{\sqrt{A} \sin \theta} \left\{ \Psi_{,\theta}, -\sqrt{\Delta} \Psi_{,r}, \frac{2I\sqrt{\Sigma}}{\alpha} \right\}$$
$$\tilde{E}(r, \theta) = \frac{\Omega - \omega}{\alpha\sqrt{\Sigma}} \left\{ \sqrt{\Delta} \Psi_{,r}, \Psi_{,\theta}, 0 \right\} .$$

$$\rho_e \tilde{E} + \tilde{J} \times \tilde{B} = 0$$

The GR pulsar equation

$$\begin{aligned}
 & \left\{ \Psi_{,rr} + \frac{1}{\Delta} \Psi_{,\theta\theta} + \Psi_{,r} \left(\frac{A_{,r}}{A} - \frac{\Sigma_{,r}}{\Sigma} \right) - \frac{\Psi_{,\theta} \cos \theta}{\Delta \sin \theta} \right\} \cdot \left[1 - \frac{\omega^2 A \sin^2 \theta}{\Sigma} + \frac{4M\alpha\omega r \sin^2 \theta}{\Sigma} - \frac{2Mr}{\Sigma} \right] \\
 & - \left(\frac{A_{,r}}{A} - \frac{\Sigma_{,r}}{\Sigma} \right) \Psi_{,r} - \left(2 \frac{\cos \theta}{\sin \theta} - \frac{A_{,\theta}}{A} + \frac{\Sigma_{,\theta}}{\Sigma} \right) (\omega^2 A \sin^2 \theta - 4M\alpha\omega r \sin^2 \theta + 2Mr) \frac{\Psi_{,\theta}}{\Delta \Sigma} \\
 & + \frac{2Mr}{\Sigma} \left(\frac{A_{,r}}{A} - \frac{1}{r} \right) \Psi_{,r} + \frac{4\omega M\alpha r \sin^2 \theta}{\Sigma} \left\{ \Psi_{,r} \left(\frac{1}{r} - \frac{A_{,r}}{A} \right) - \frac{\Psi_{,\theta} A_{,\theta}}{\Delta A} \right\} \\
 & - \frac{\omega' \sin^2 \theta}{\Sigma} (\omega A - 2\alpha M r) \left(\Psi_{,r}^2 + \frac{1}{\Delta} \Psi_{,\theta}^2 \right) = -\frac{4\Sigma}{\Delta} II'
 \end{aligned}$$

$$1 - \frac{\omega^2 A \sin^2 \theta}{\Sigma} + \frac{4M\alpha\omega r \sin^2 \theta}{\Sigma} - \frac{2Mr}{\Sigma} = 0$$

The GR pulsar equation

$$\begin{aligned}
 & \left\{ \Psi_{,rr} + \frac{1}{\Delta} \Psi_{,\theta\theta} + \Psi_{,r} \left(\frac{A_{,r}}{A} - \frac{\Sigma_{,r}}{\Sigma} \right) - \frac{\Psi_{,\theta} \cos \theta}{\Delta \sin \theta} \right\} \cdot \left[1 - \frac{\omega^2 A \sin^2 \theta}{\Sigma} + \frac{4M\alpha\omega r \sin^2 \theta}{\Sigma} - \frac{2Mr}{\Sigma} \right] \\
 & - \left(\frac{A_{,r}}{A} - \frac{\Sigma_{,r}}{\Sigma} \right) \Psi_{,r} - \left(2 \frac{\cos \theta}{\sin \theta} - \frac{A_{,\theta}}{A} + \frac{\Sigma_{,\theta}}{\Sigma} \right) (\omega^2 A \sin^2 \theta - 4M\alpha\omega r \sin^2 \theta + 2Mr) \frac{\Psi_{,\theta}}{\Delta \Sigma} \\
 & + \frac{2Mr}{\Sigma} \left(\frac{A_{,r}}{A} - \frac{1}{r} \right) \Psi_{,r} + \frac{4\omega M\alpha r \sin^2 \theta}{\Sigma} \left\{ \Psi_{,r} \left(\frac{1}{r} - \frac{A_{,r}}{A} \right) - \frac{\Psi_{,\theta} A_{,\theta}}{\Delta A} \right\} \\
 & - \frac{\omega' \sin^2 \theta}{\Sigma} (\omega A - 2\alpha M r) \left(\Psi_{,r}^2 + \frac{1}{\Delta} \Psi_{,\theta}^2 \right) = -\frac{4\Sigma}{\Delta} II'
 \end{aligned}$$

$$1 - \frac{\omega^2 A \sin^2 \theta}{\Sigma} + \frac{4M\alpha\omega r \sin^2 \theta}{\Sigma} - \frac{2Mr}{\Sigma} = 0$$

The GR pulsar equation

$$\left(\Psi_{,rr} + \frac{1}{r^2} \Psi_{,\theta\theta} + \frac{2\Psi_{,r}}{r} - \frac{1}{r^2} \frac{\cos \theta}{\sin \theta} \Psi_{,\theta} \right) \cdot [1 - \omega^2 r^2 \sin^2 \theta]$$
$$- \frac{2\Psi_{,r}}{r} - 2\omega^2 \cos \theta \sin \theta \Psi_{,\theta} - \omega\omega' r^2 \sin^2 \theta \left(\Psi_{,r}^2 + \frac{1}{r^2} \Psi_{,\theta}^2 \right) = -4II'$$

- The pulsar light cylinder: $r \sin \theta = c / \omega$
- The electric current $I(\Psi)$ must be determined self-consistently

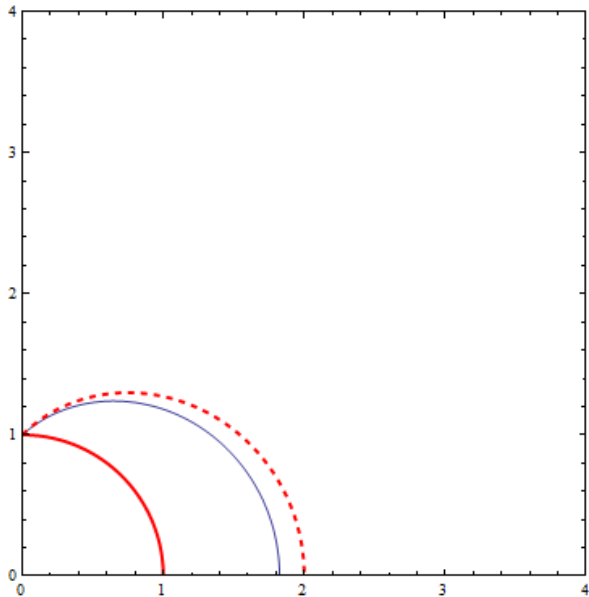
The GR pulsar equation

$$1 - \frac{\omega^2 A \sin^2 \theta}{\Sigma} + \frac{4M\alpha\omega r \sin^2 \theta}{\Sigma} - \frac{2Mr}{\Sigma} = 0$$

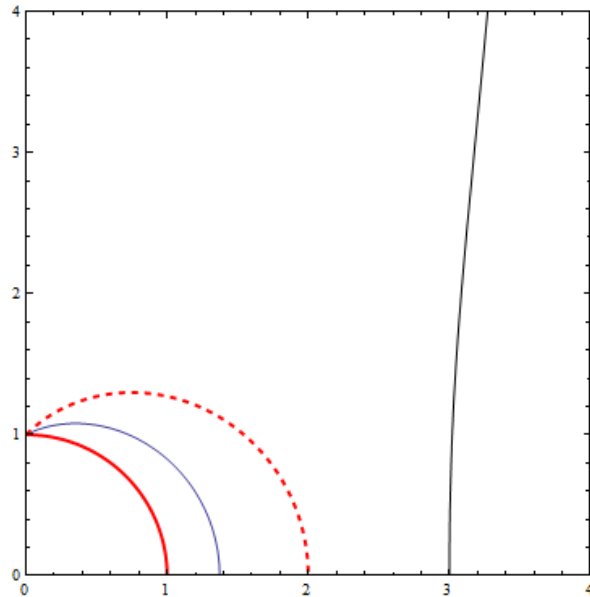
- The black hole possesses **two light surfaces**
- The electric current $I(\Psi)$ must be determined self-consistently together with the angular velocity of the magnetic field $\omega(\Psi)$

The GR pulsar equation

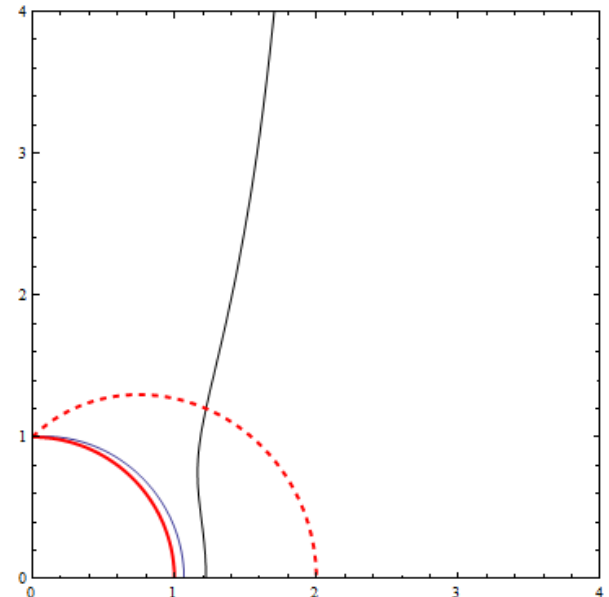
$\alpha=M, \omega=0.1 \Omega_{\text{BH}}$



$\alpha=M, \omega=0.5 \Omega_{\text{BH}}$

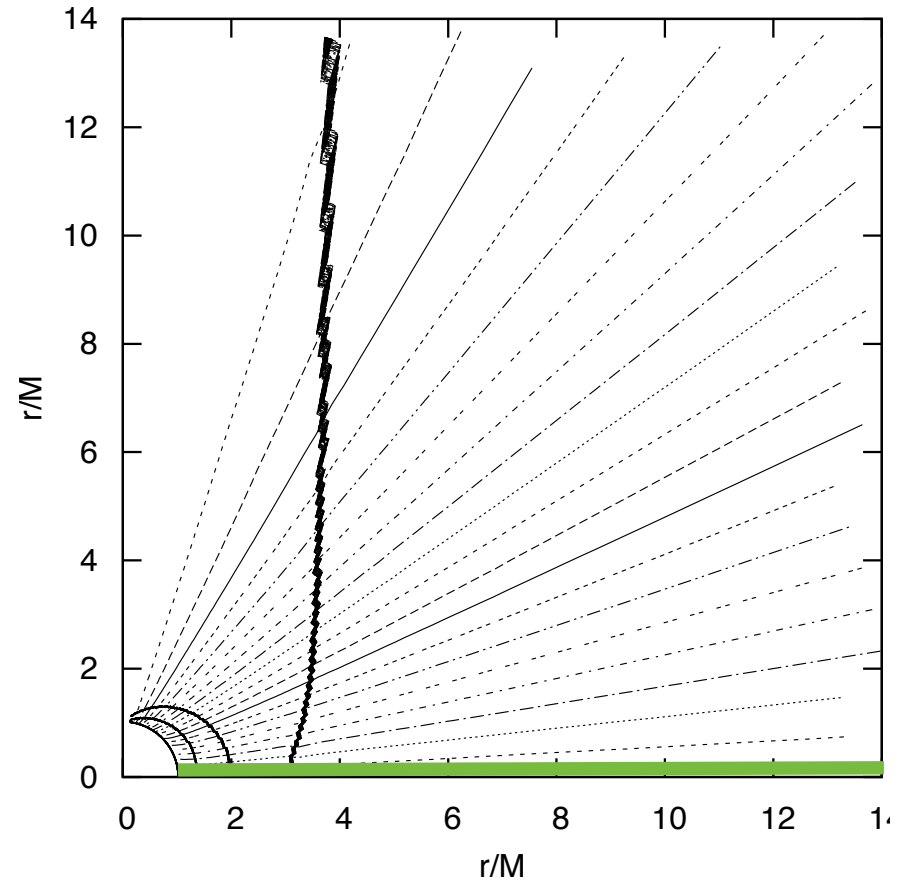
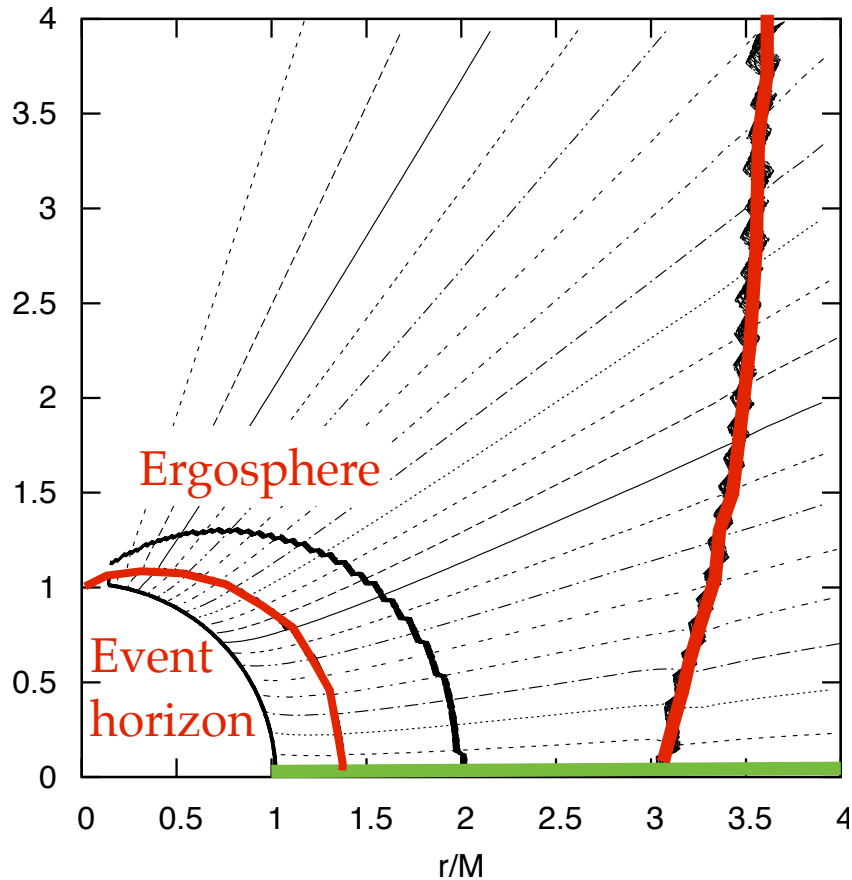


$\alpha=M, \omega=0.9 \Omega_{\text{BH}}$



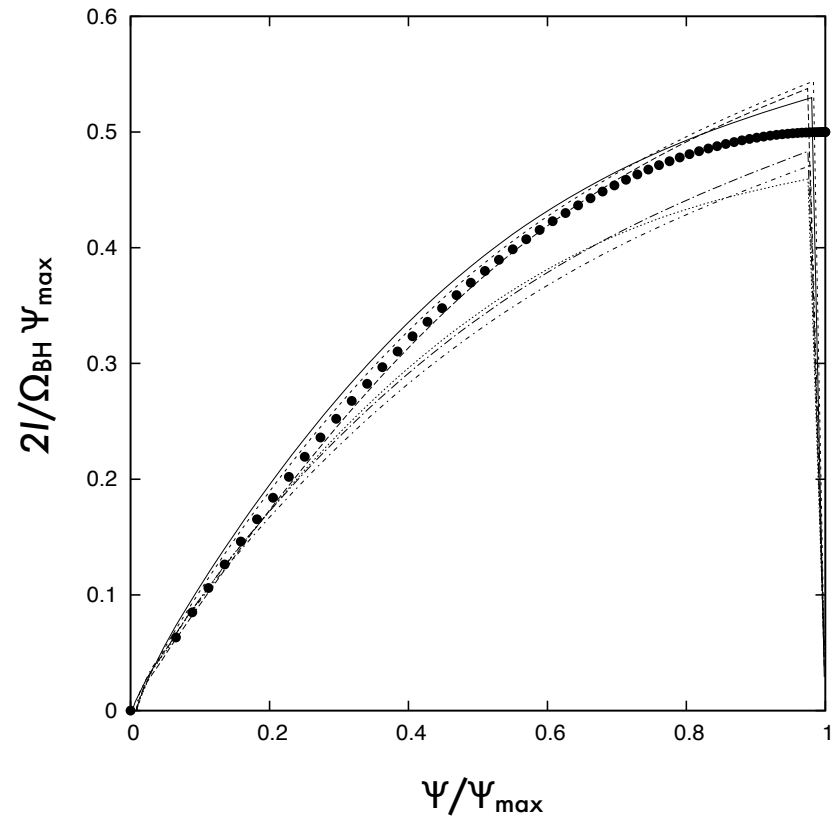
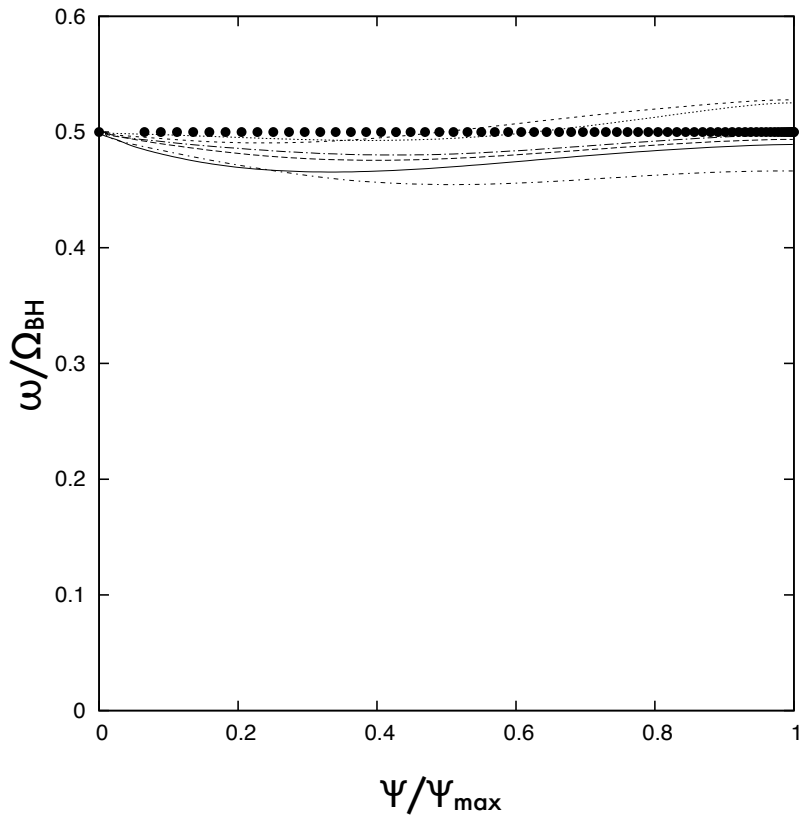
The GR pulsar equation

$$\alpha=0.9999M, \omega \sim 0.5 \Omega_{\text{BH}}$$

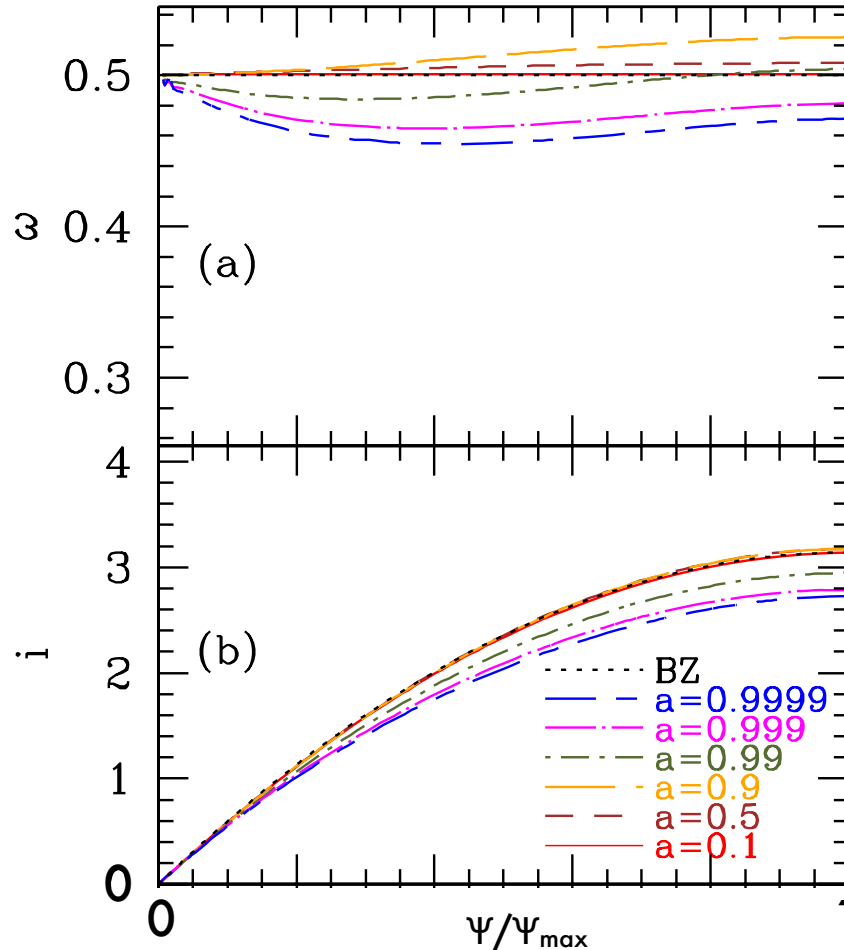


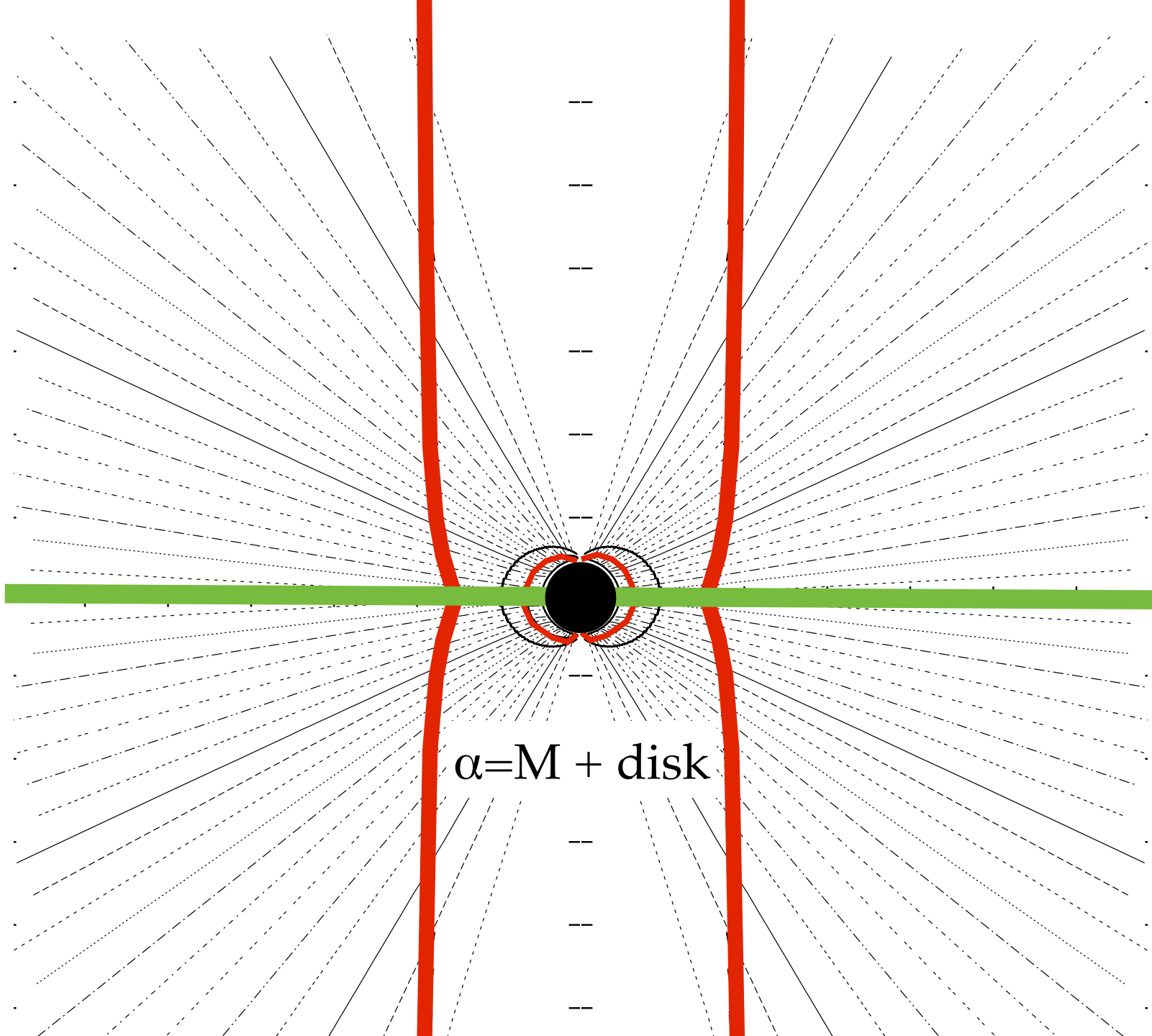
Blandford-Znajek revisited

$$\alpha=0.7-0.9999M, \omega \sim 0.5 \Omega_{\text{BH}}$$

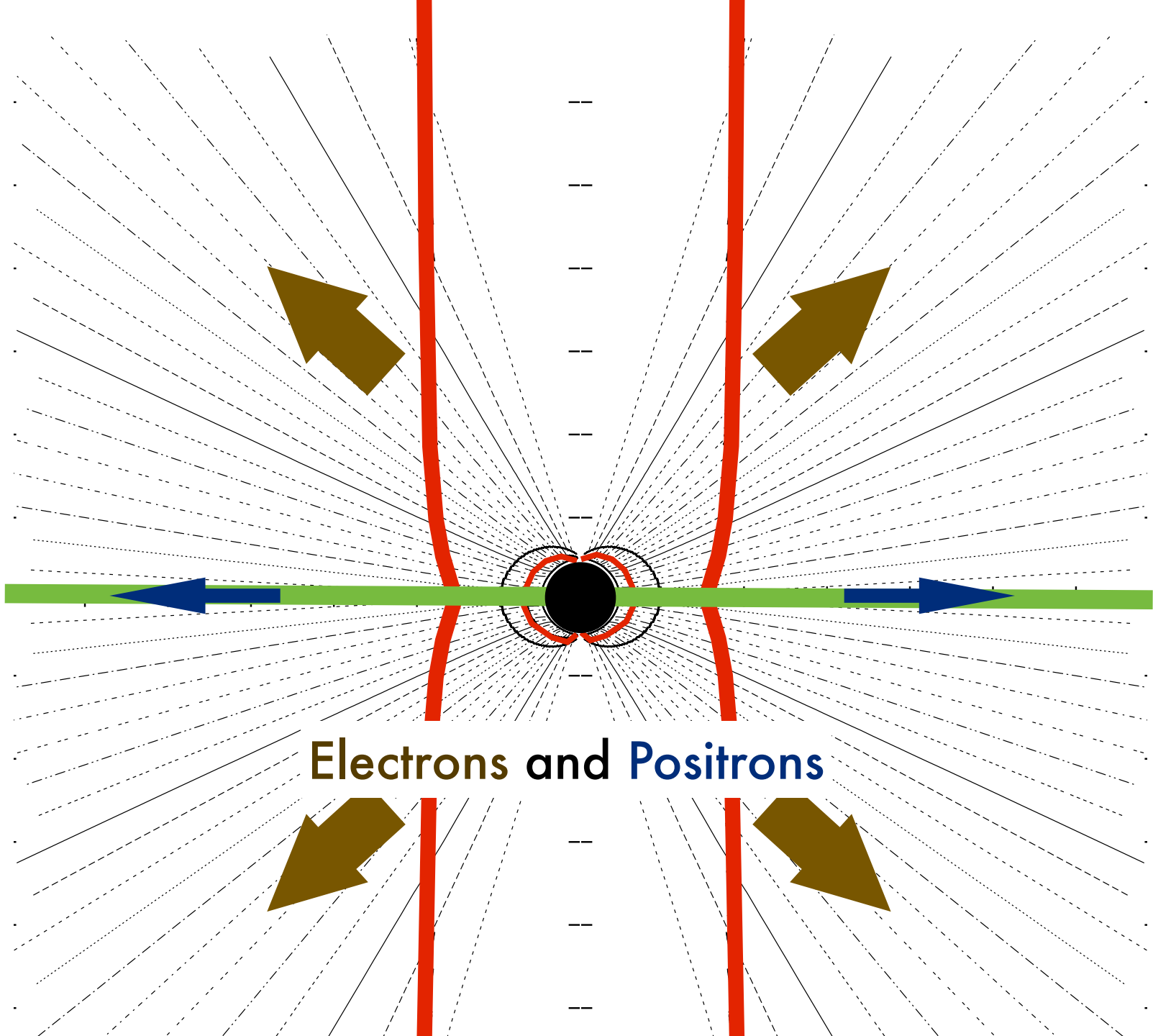


Blandford-Znajek revisited

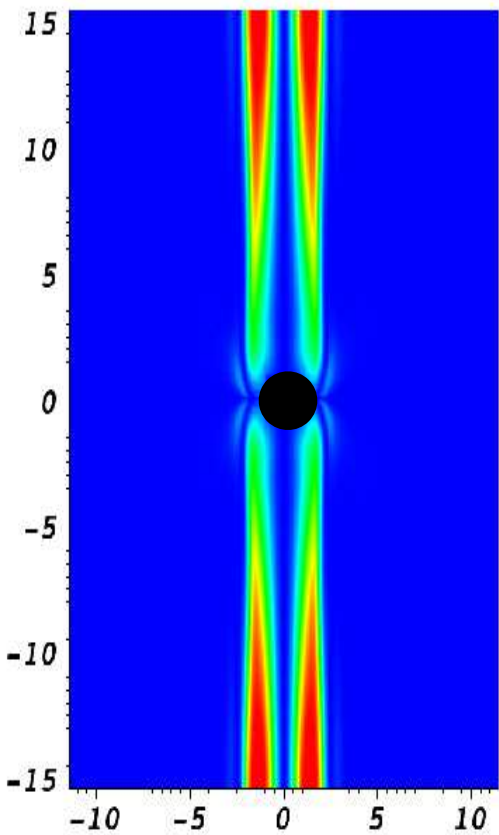




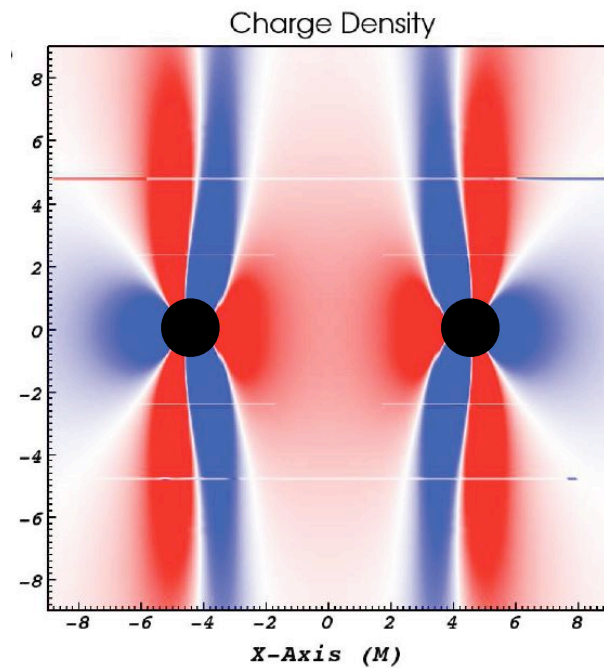
$\alpha = M + \text{disk}$



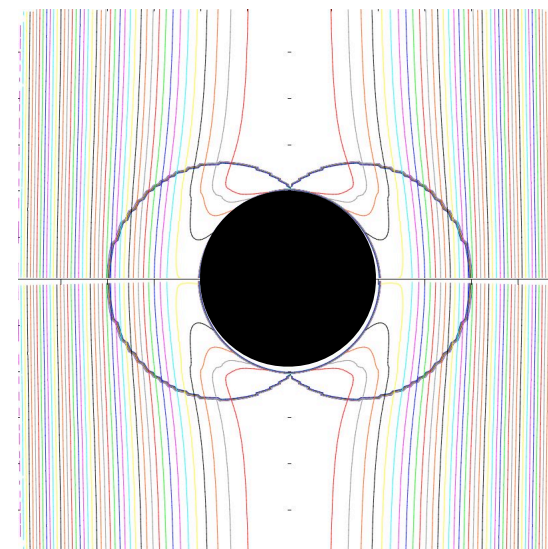
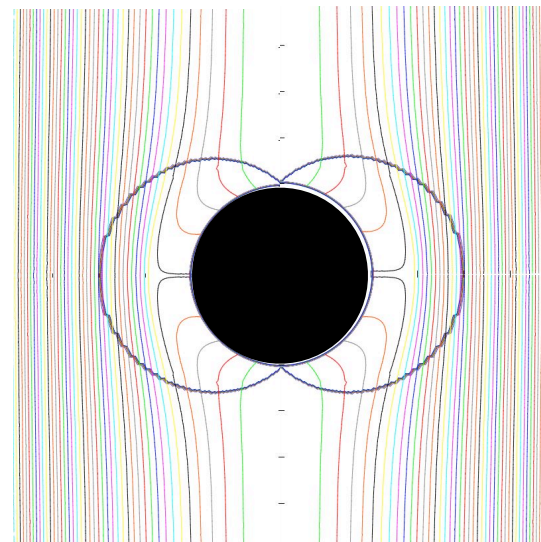
Electrons and Positrons



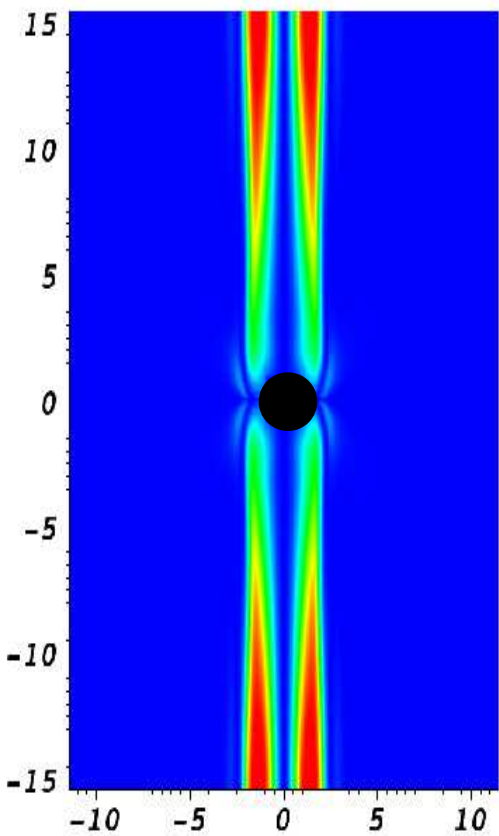
Palenzuela, Bona,
Lehner, Reula 2011



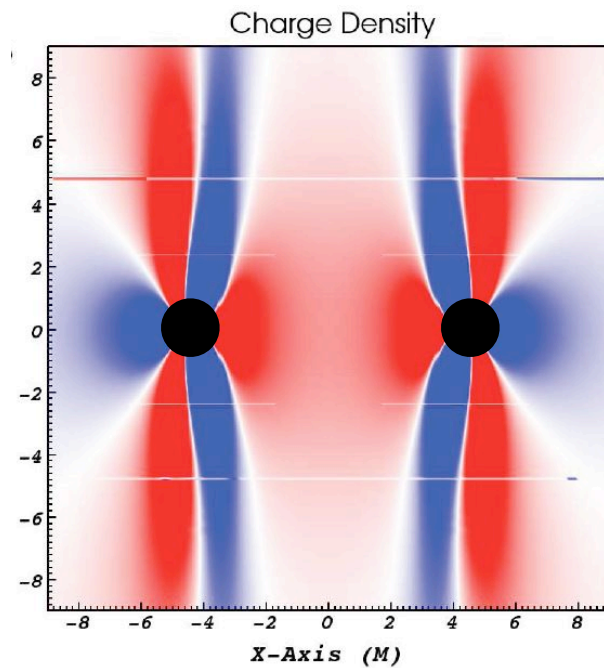
Alic, Moesta, Rezzolla,
Jaramillo, Palenzuela,
Zanotti 2013



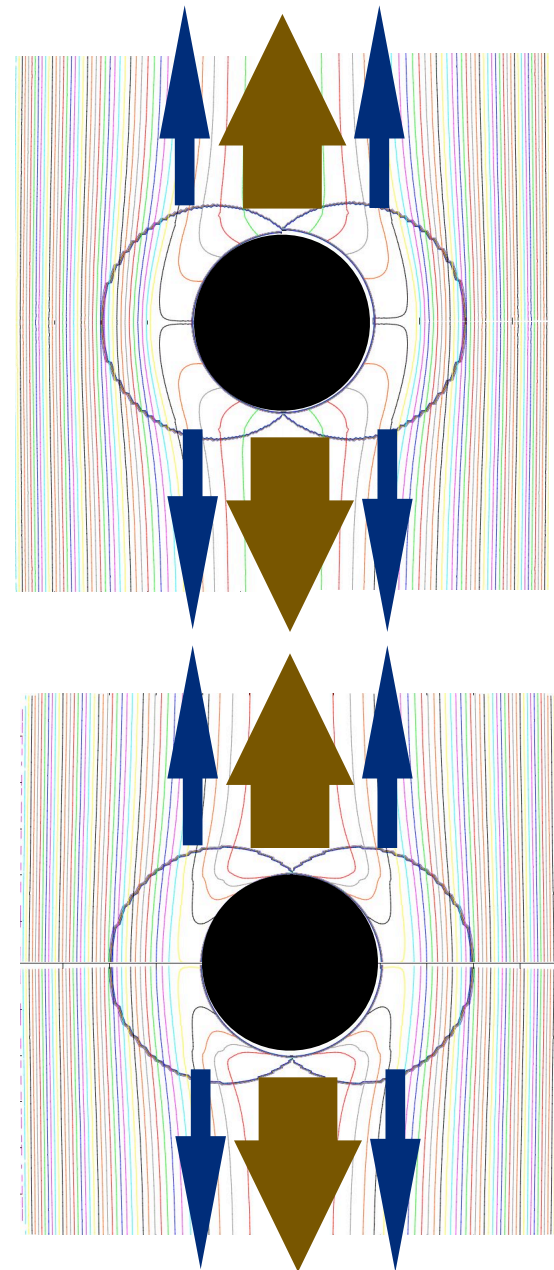
Contopoulos et al. 2013
(in preparation)



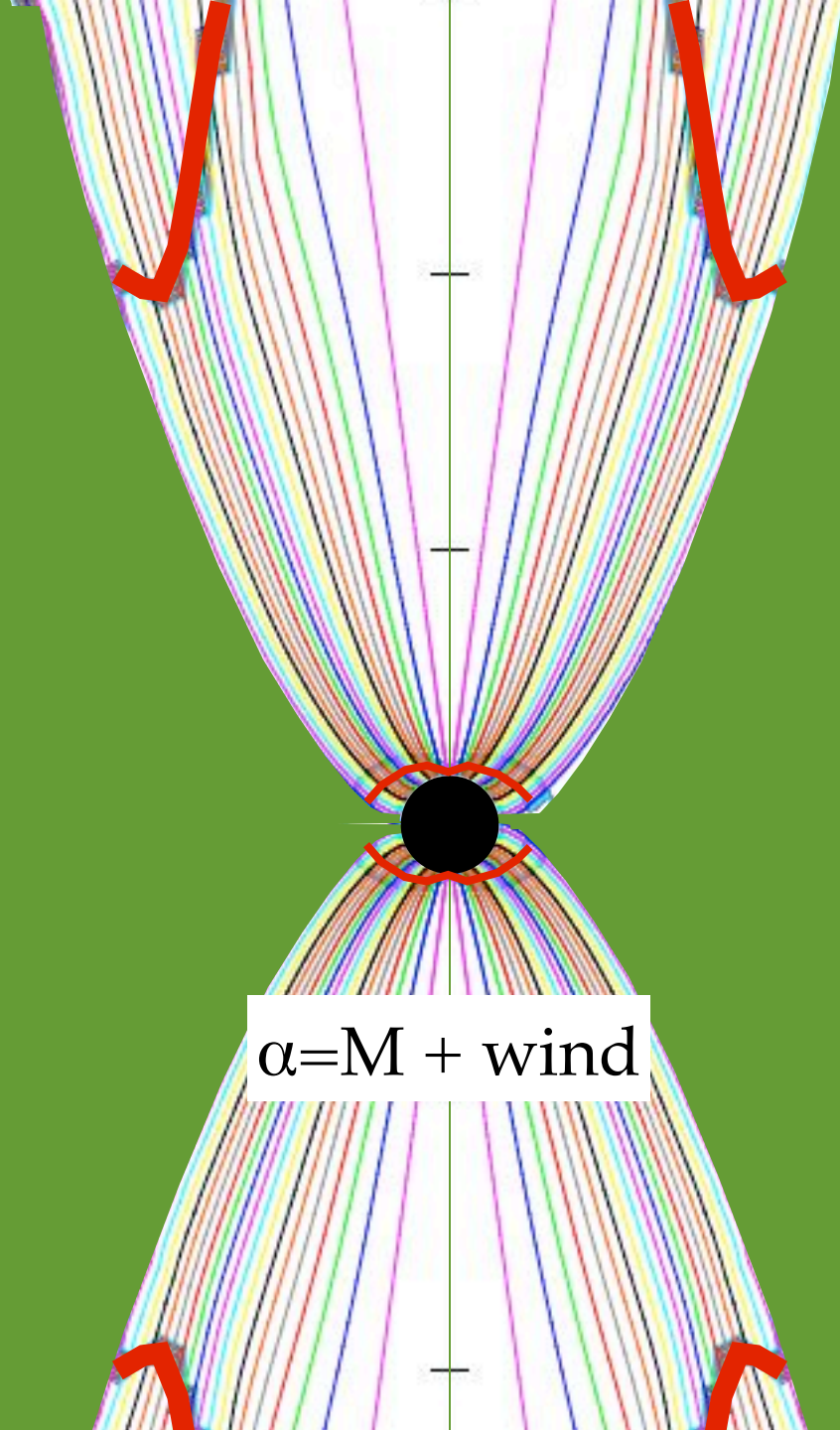
Palenzuela, Bona,
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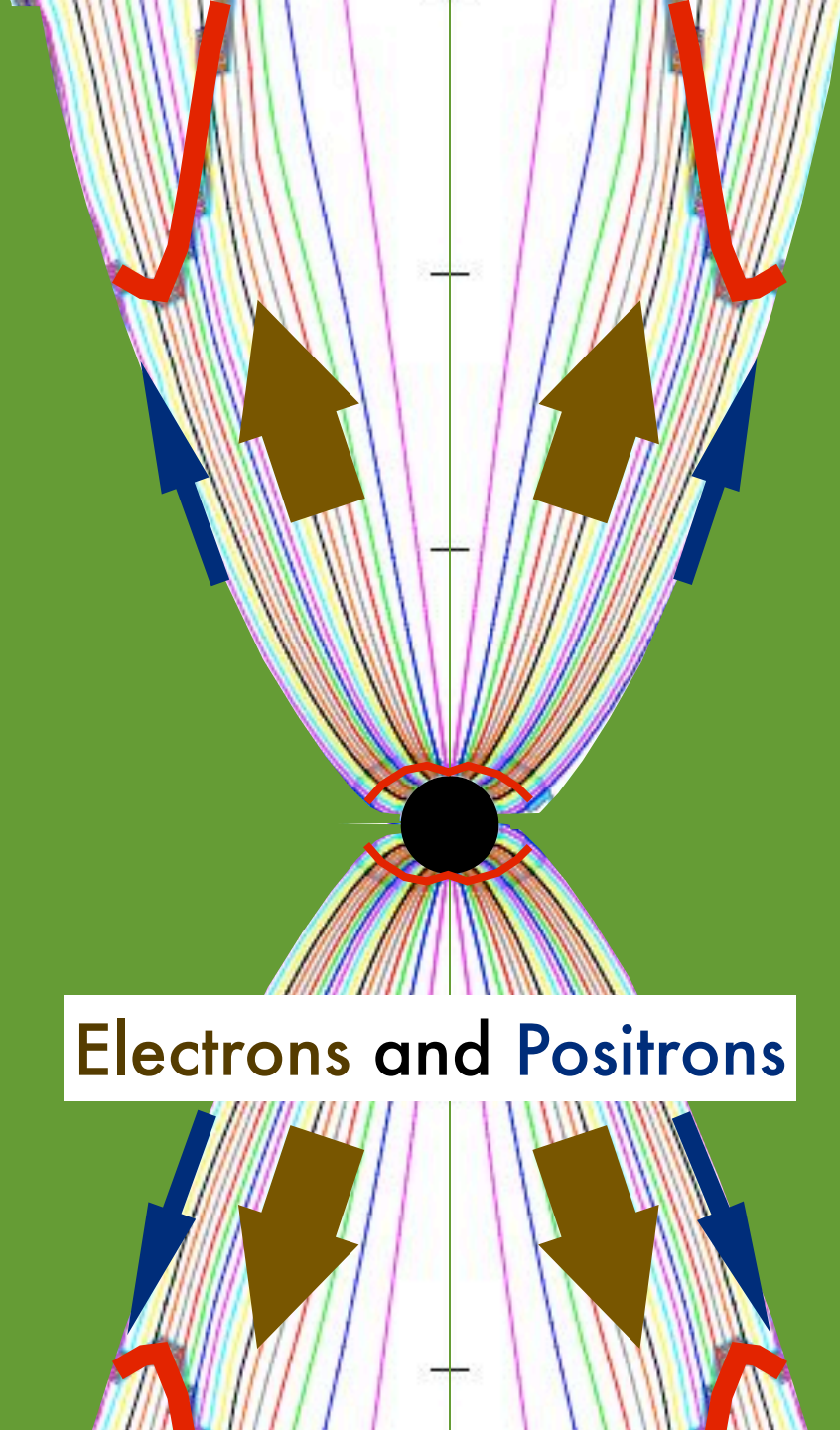
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Contopoulos et al. 2013
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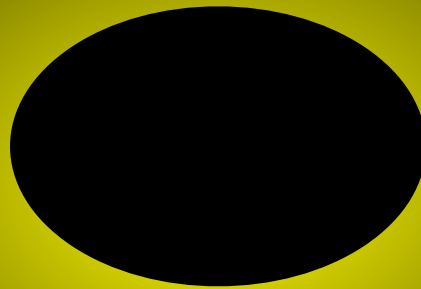


$$\alpha = M + \text{wind}$$

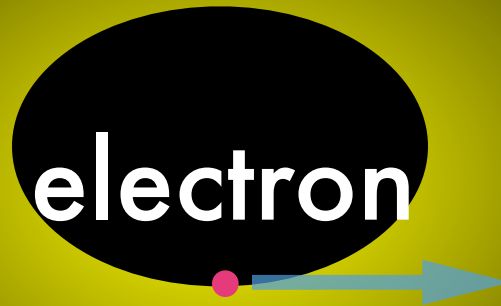


Electrons and Positrons

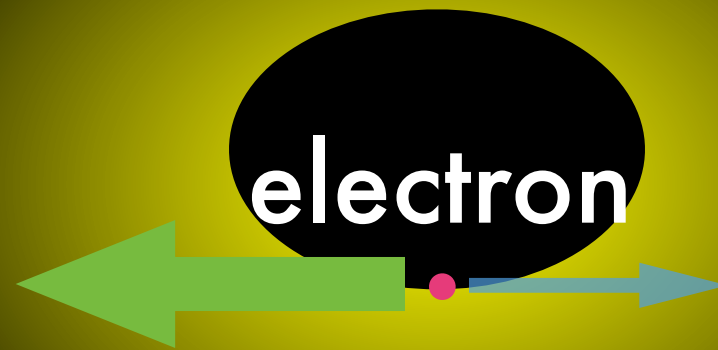
The Cosmic Battery



The Cosmic Battery



The Cosmic Battery

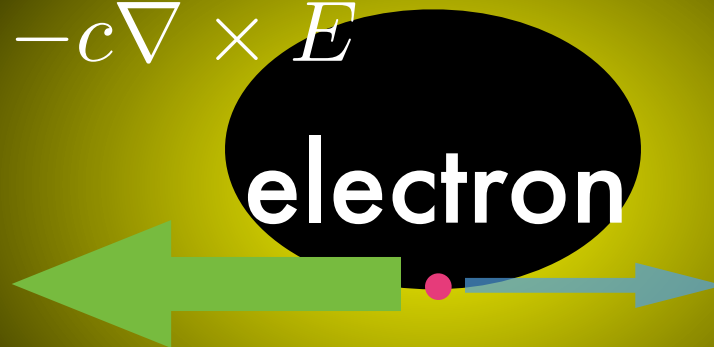


$$F_{PR} = - \frac{L \sigma_T}{4\pi r^2 c} \frac{v_\phi}{c}$$

The Cosmic Battery

$$\nabla \times B = \frac{4\pi}{c} J \quad -eE + F_{PR} = m_e \frac{dv_e}{dt}$$

$$\frac{\partial B}{\partial t} = -c \nabla \times E \quad eE = m_p \frac{dv_p}{dt}$$



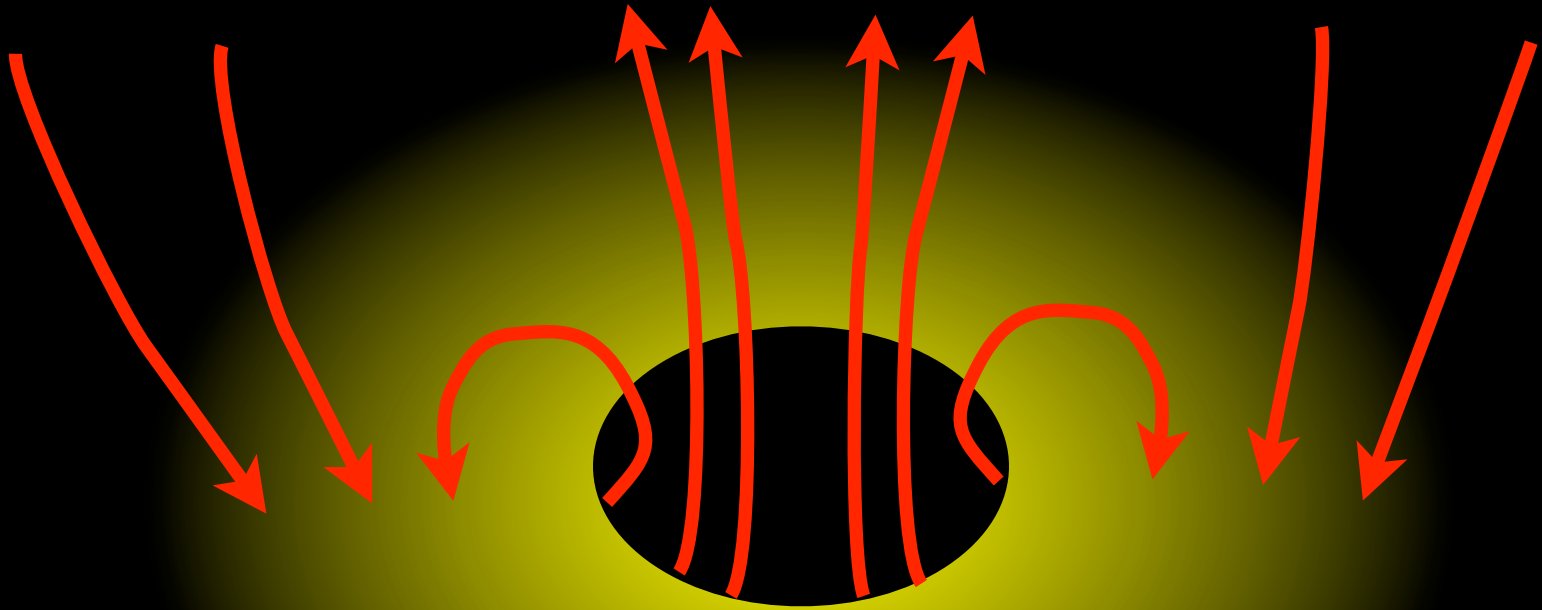
$$F_{PR} = -\frac{L\sigma_T}{4\pi r^2 c} \frac{v_\phi}{c}$$

The Cosmic Battery



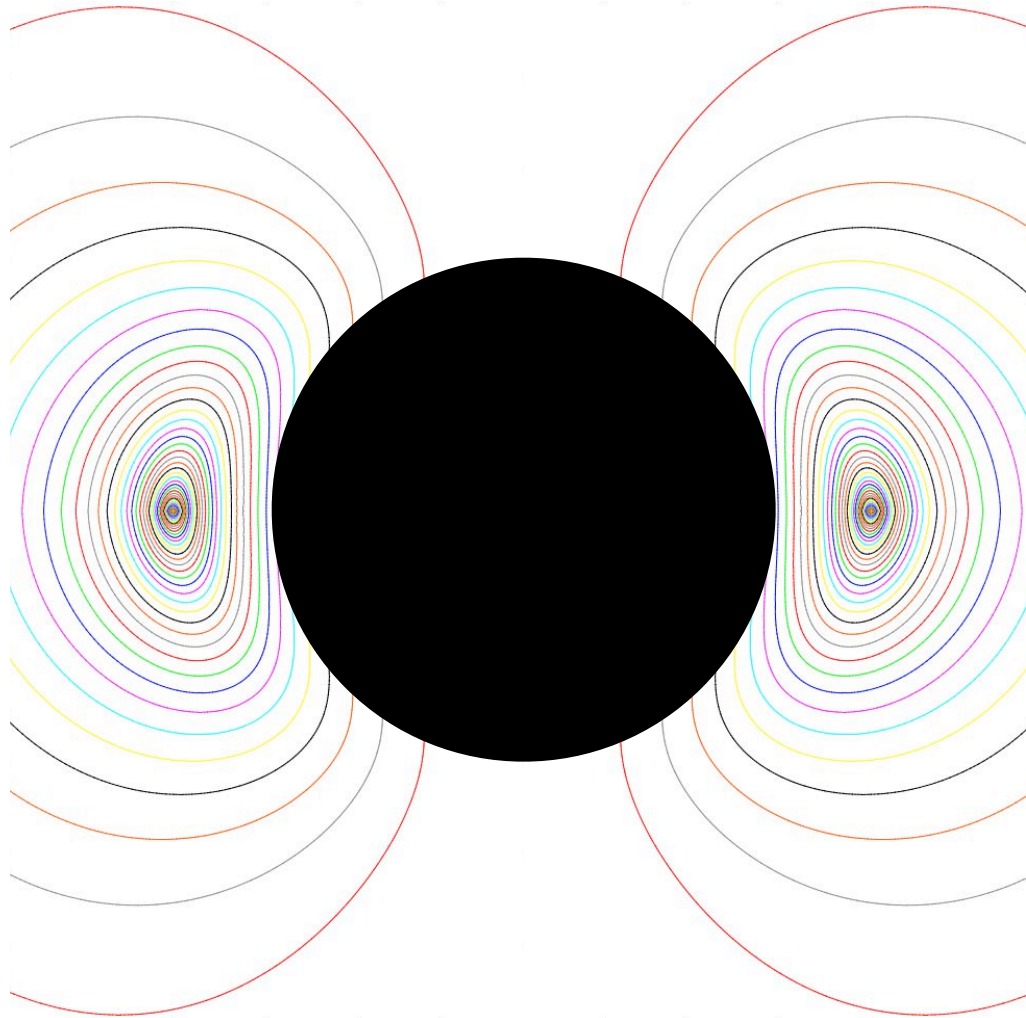
$$\frac{\partial B}{\partial t} = -c \nabla \times (\dots + F_{PR}/e)$$

The Cosmic Battery



$$\frac{\partial B}{\partial t} = -c \nabla \times (\dots + F_{PR}/e)$$

The Cosmic Battery



Conclusions

- In analogy with pulsars...
 - The light cylinder determines a unique solution
 - Isolated black holes do not produce jets
 - “Operational” black holes need very efficient pair formation above the horizon
- The background magnetic field may be **generated in situ by the Cosmic Battery**