

Broadband radio circular polarization spectrum of the relativistic jet in PKS B2126-158

O'Sullivan, McClure-Griffiths, Feain, Gaensler, Sault (2013) MNRAS, in review

Shane O'Sullivan, Relativistic Jets & Their Magnetic Fields, Spain, 13 Jun 2013



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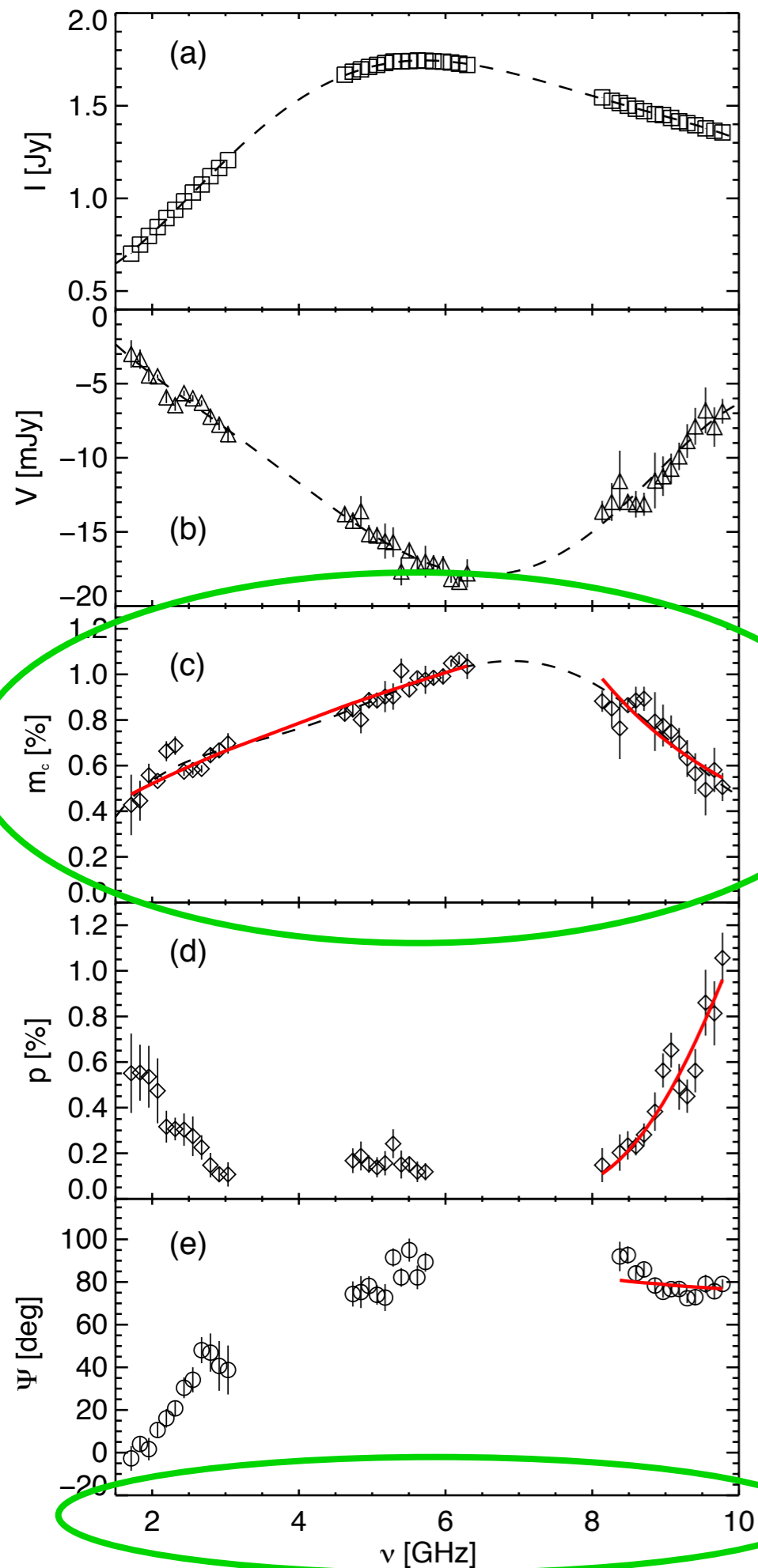
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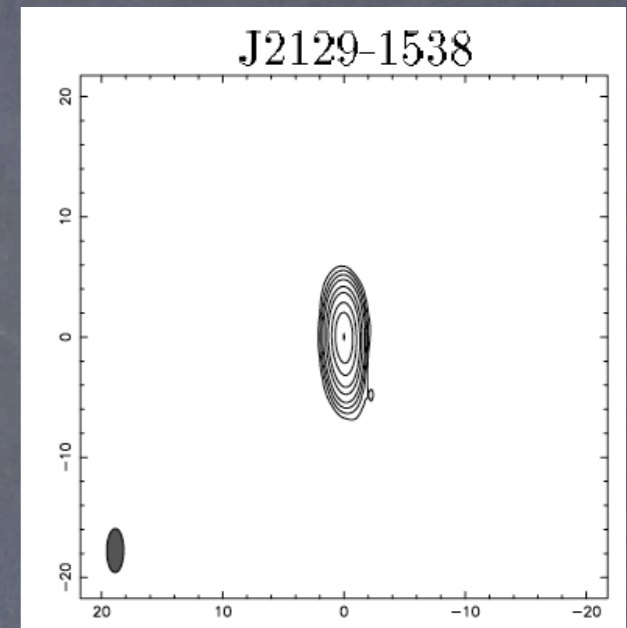
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- GPS quasar
- $z = 3.268$
- $M_{\text{BH}} \sim 10^{10} M_{\text{sol}}$
- $P_{\text{jet}} \sim 10^{48} \text{ erg/s}$
- $m_c \sim 1\% \text{ at } 6 \text{ GHz}$



- **Dominant CP generation mechanism: Faraday conversion of LP to CP**

- Conclusions:

1. little or no thermal plasma within the jet
2. cannot yet distinguish between magnetic twist and internal Faraday rotation models
3. relativistic particle content unconstrained

Australia Telescope Compact Array

- 6 x 22 m dishes with 6 km maximum baseline
- Linear feed system ($s_{V/I} \sim 0.01\%$)
- CABB (2 GHz b/w per IF), Wilson et al. (2011)
- 1 to 3 GHz, 4.5 to 6.5 GHz, 8 to 10 GHz (1 to 10 arcsec)

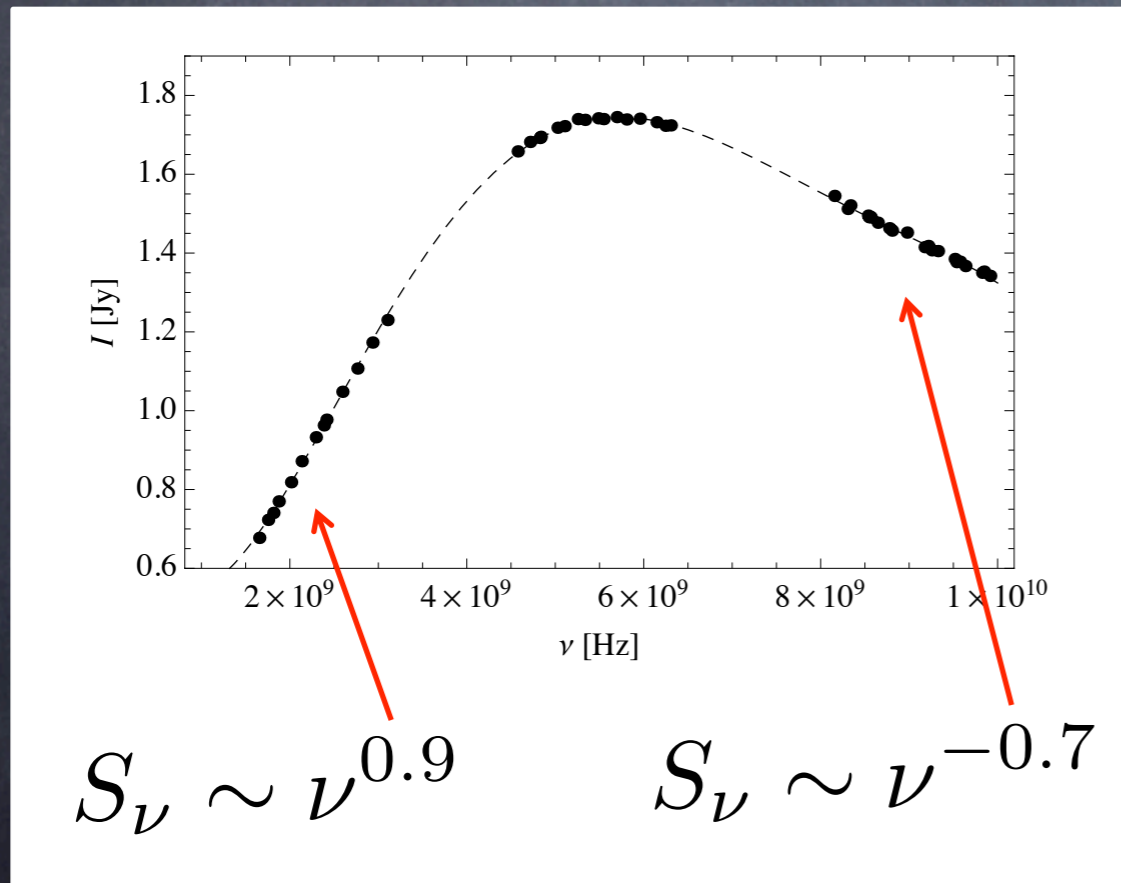
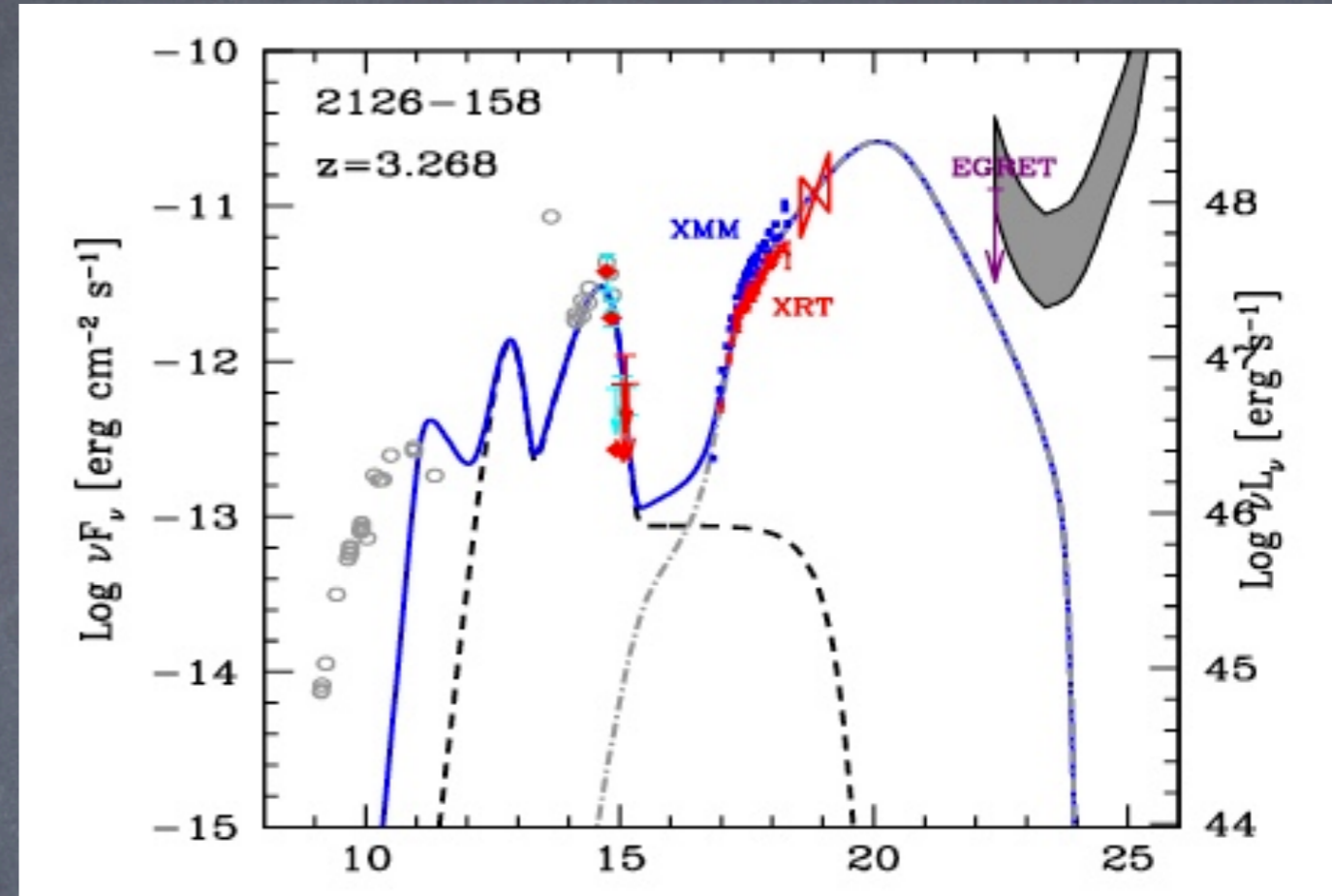
$$\begin{aligned} I &= v_{xx}^* + v_{yy}^* & LP &= \sqrt{Q^2 + U^2} \\ Q &= v_{xx}^* - v_{yy}^* & \chi &= \frac{1}{2} \tan^{-1}(U/Q) \\ U &= v_{xy}^* + v_{yx}^* & CP &= V \\ V &= v_{xy}^* - v_{yx}^* \end{aligned}$$



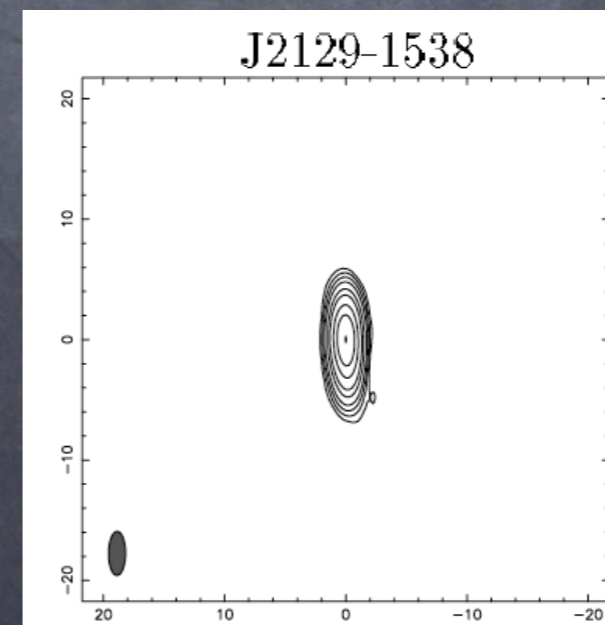
PKS B2126-158

- GPS quasar (Stangellini+98)
- $z = 3.268$ (Osmer+94)
- $M_{\text{BH}} \sim 10^{10} M_{\text{sol}}$ (Ghisellini+11)

Ghisellini+11



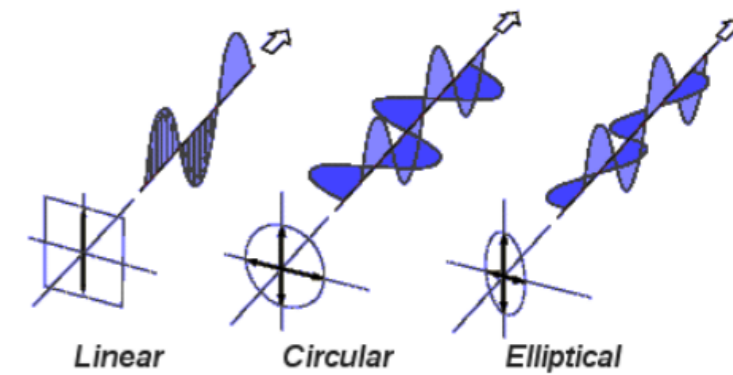
VLBA 5 GHz



Fomalont+95

- Synchrotron from a single electron is elliptically polarized
- CP can be positive/negative (left/right)
- Large CP not expected for an ensemble of electrons with an isotropic distribution of pitch angles

Polarization of electromagnetic waves



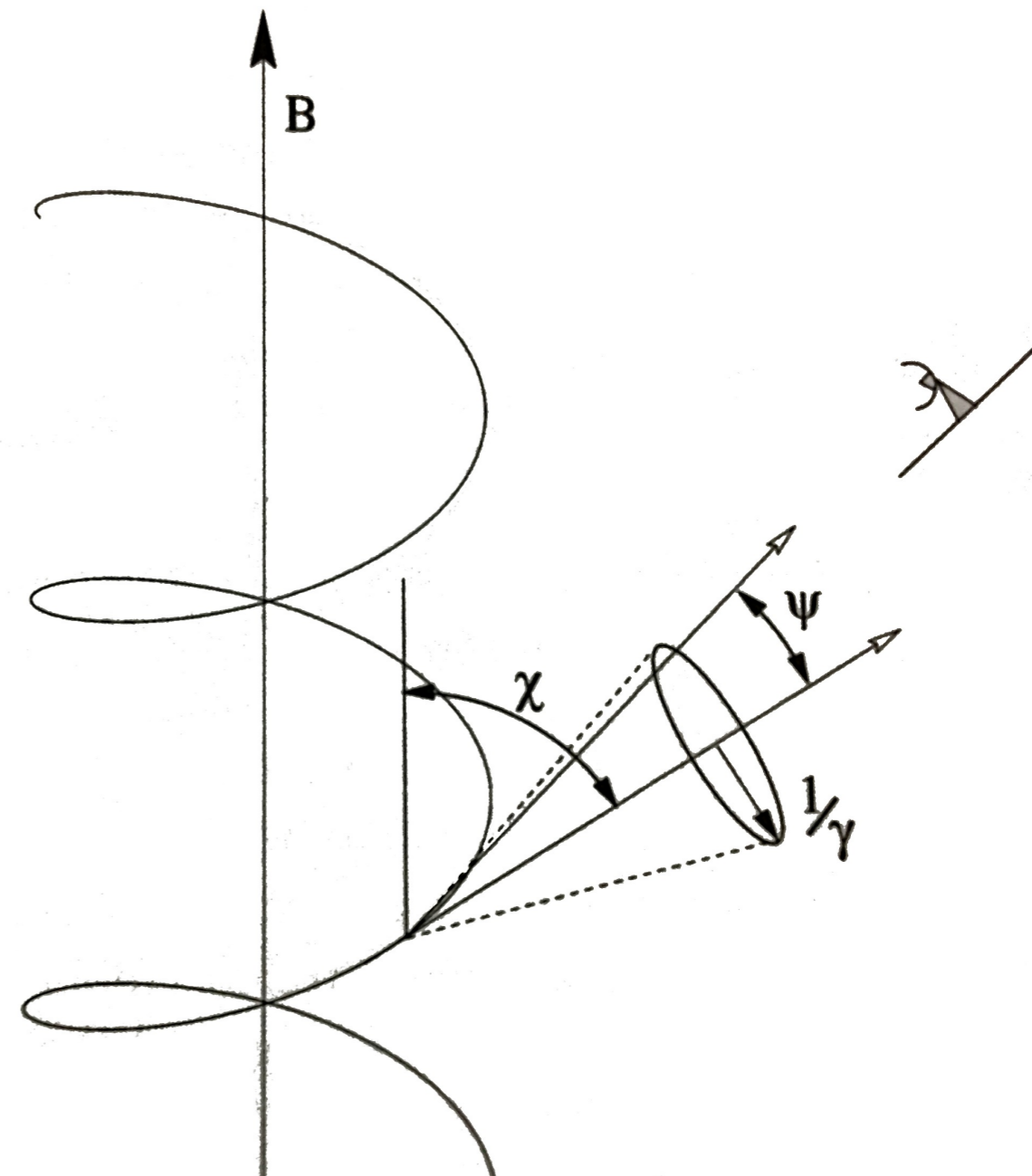
- direct measure of magnetic field strength

$$m_c \sim \left(\frac{2.8B[G] \sin \theta}{\nu[MHz]} \right)^{1/2} \cot \theta$$

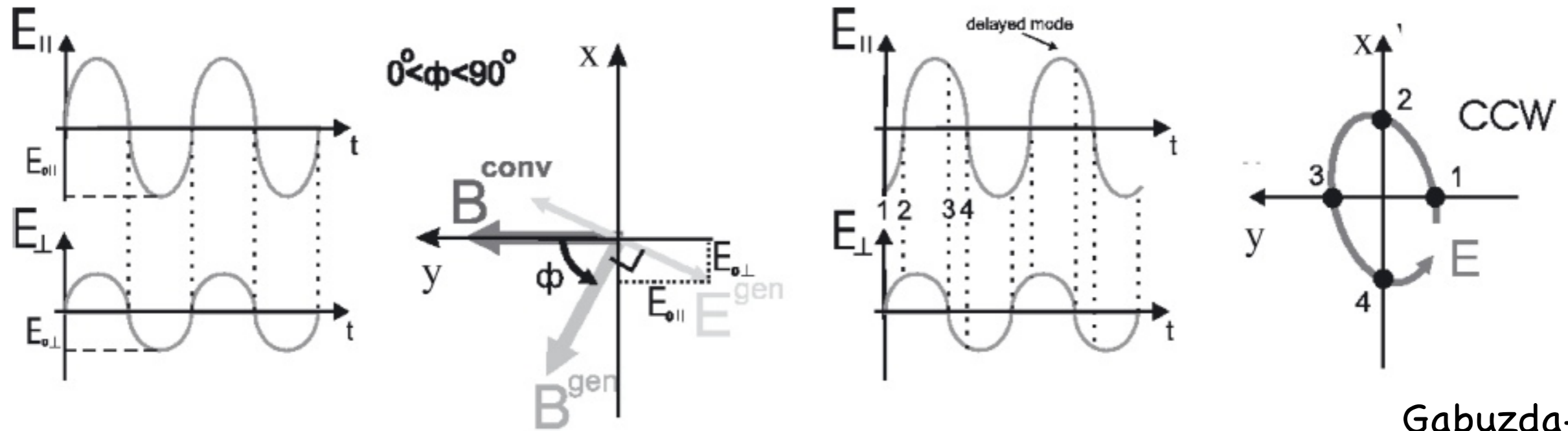
- 10 mG, 1 GHz (1 G, 100 GHz) -> ~0.3%

Legg & Westfold (1968)

- For completely uniform magnetic field with an ensemble of electrons and random distribution of pitch angles gives LP ~ 75%
- Polarization => strength and degree of order of magnetic field

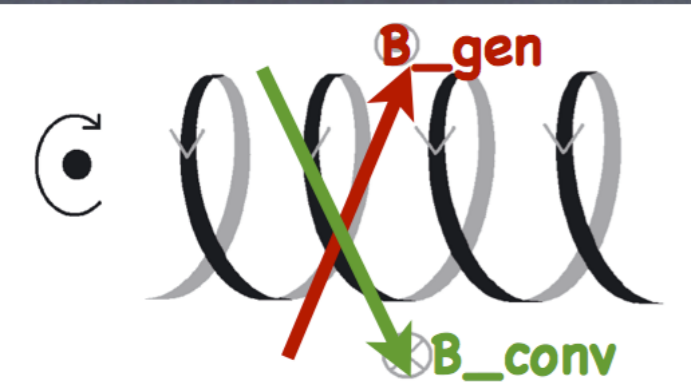


Conversion of LP to CP

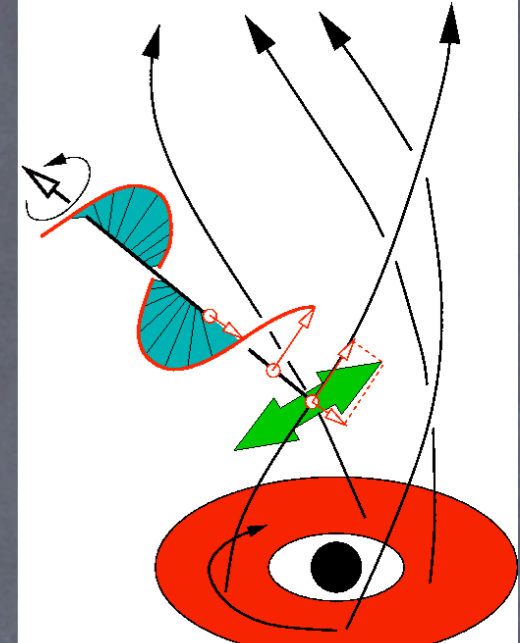


Gabuzda+08

- Propagation effect due to finite width of the jet (B^{conv} induces a delay between E_{\parallel} and E_{\perp})
- A **change in the orientation of the magnetic field** through the jet eg. helical magnetic field
- Conversion due to **Faraday rotation**: Linear polarization emitted from back of the jet rotated as it propagates through the jet



CP Spectrum



- Determine the CP generation mechanism
- Using correct model allows us to infer:
 - B-field geometry (& strength, flux) (Gabuzda+08, Homan+09)
 - Jet particle composition (e^- , p / e^- , e^+) (Wardle+98)
 - Direction of rotation of accretion disk/black hole spin (Ensslin+03)
 - Polarity of black hole magnetic field (Beckert & Falcke 2002)

Circular Polarization in AGN

Predicted spectrum: $m_c \equiv |V|/I$

- Intrinsic: $m_c \propto \nu^{-1/2}$

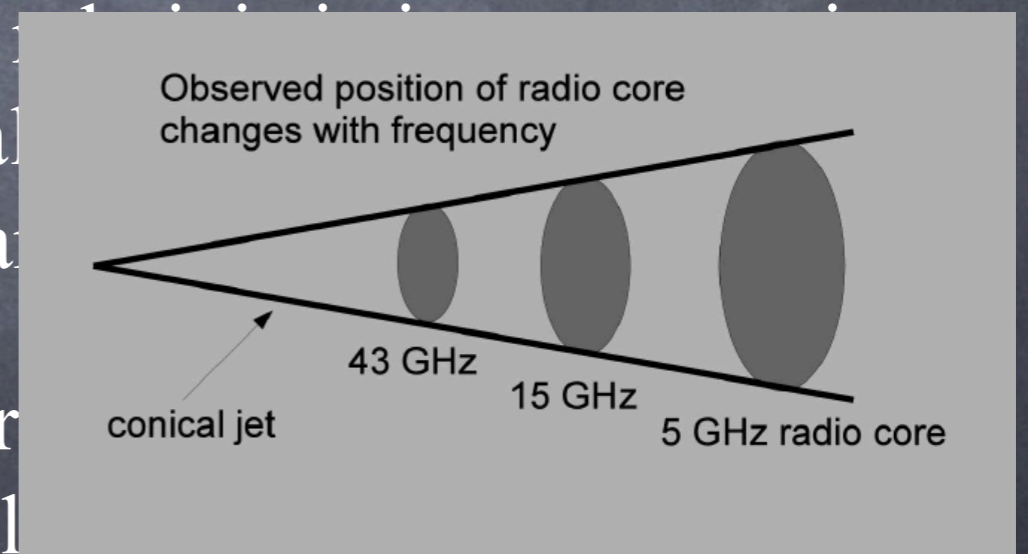
- Melrose (1972): sign change $\sim 1/2$ self-abs turnover freq

- Conversion:

- Jones (1988): inhomogeneous
greatest is region where optical
will occur regardless of mecha

- Kennett & Melrose (1998): for
cold plasma: $m_c \propto \nu^{-1}$ purely

- Wardle & Homan (2003): BK jet: $m_c \propto \nu^{+1}$



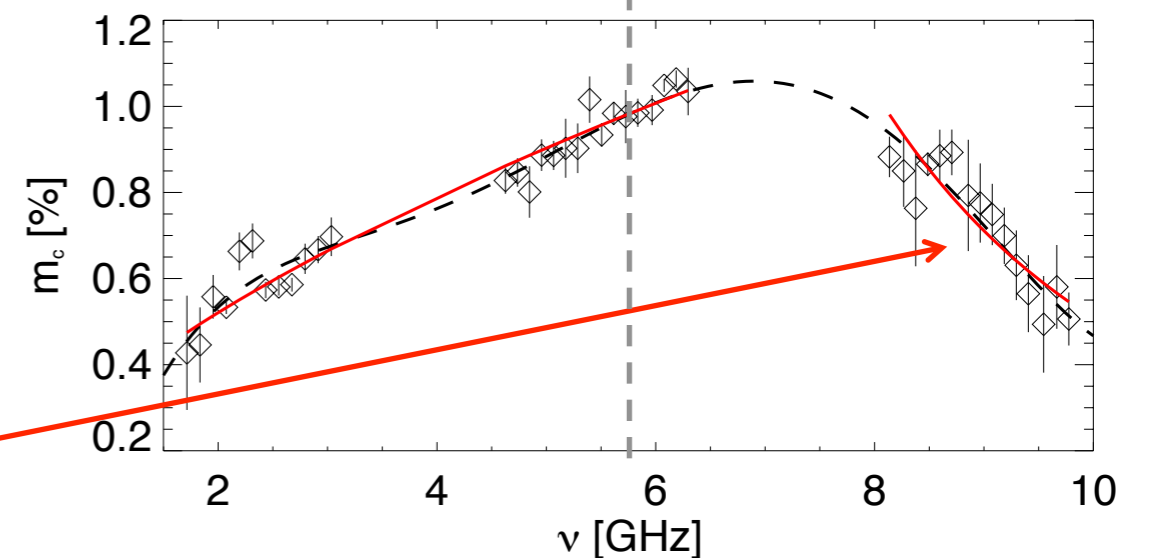
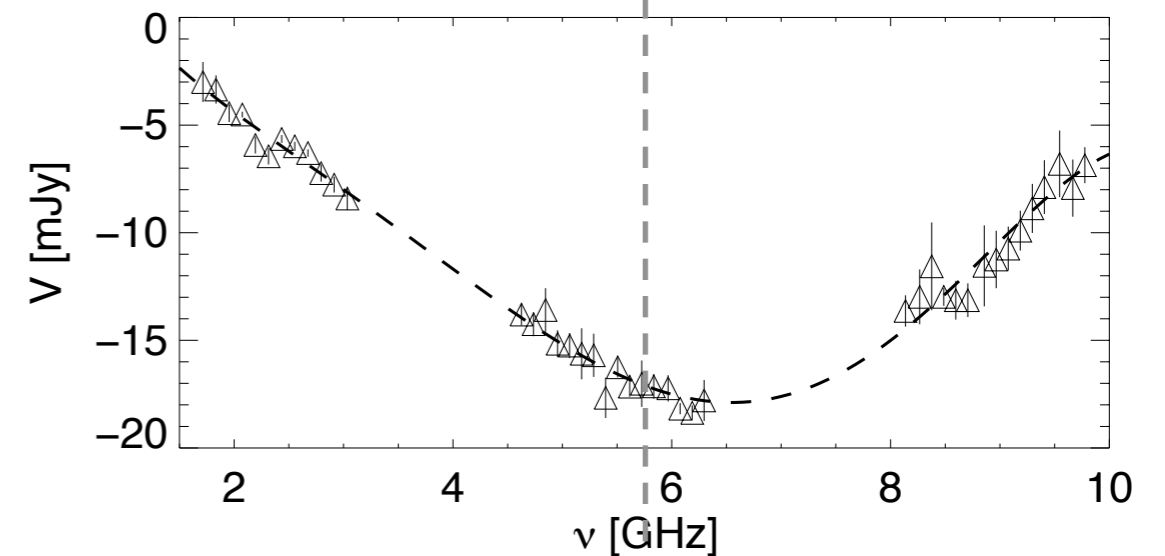
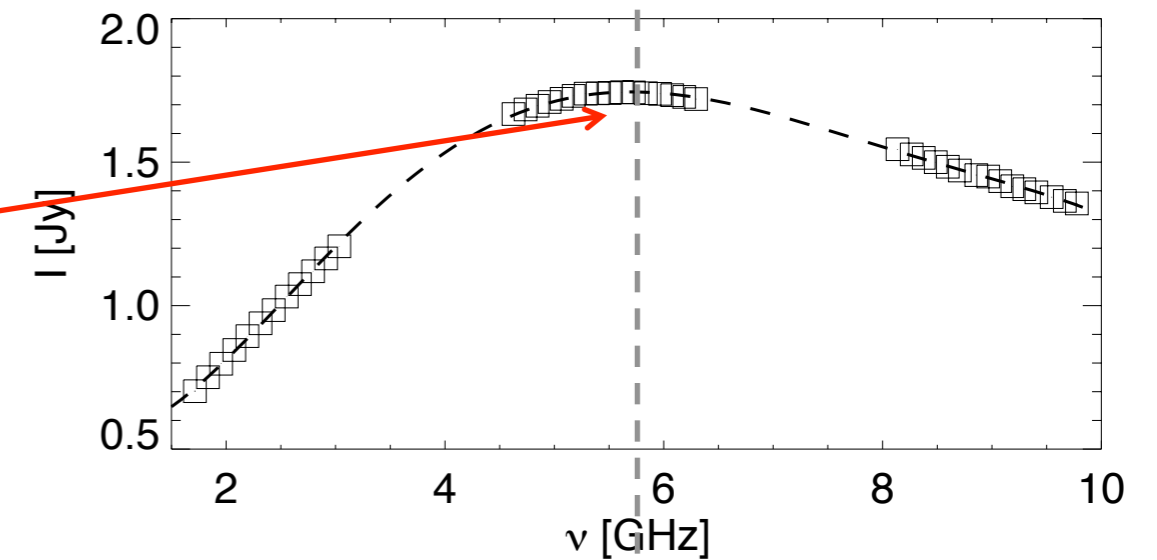
PKS B2126-128

$$\nu_t \sim 5.7 \pm 0.1 \text{ GHz}$$

$$S_{\nu_t} \sim 1.744 \pm 0.001 \text{ GHz}$$

- $m_c \sim 1\%$ at 6 GHz
- No sign change in Stokes V
- Turnover in CP spectrum when emission becomes optically thin
- Not dominated by intrinsic CP
- Consistent with Faraday conversion due to uniform B-field in a purely relativistic plasma (no thermal electrons)

$$m_c \propto \nu^{-3.0 \pm 0.4}$$



PKS B2126-128

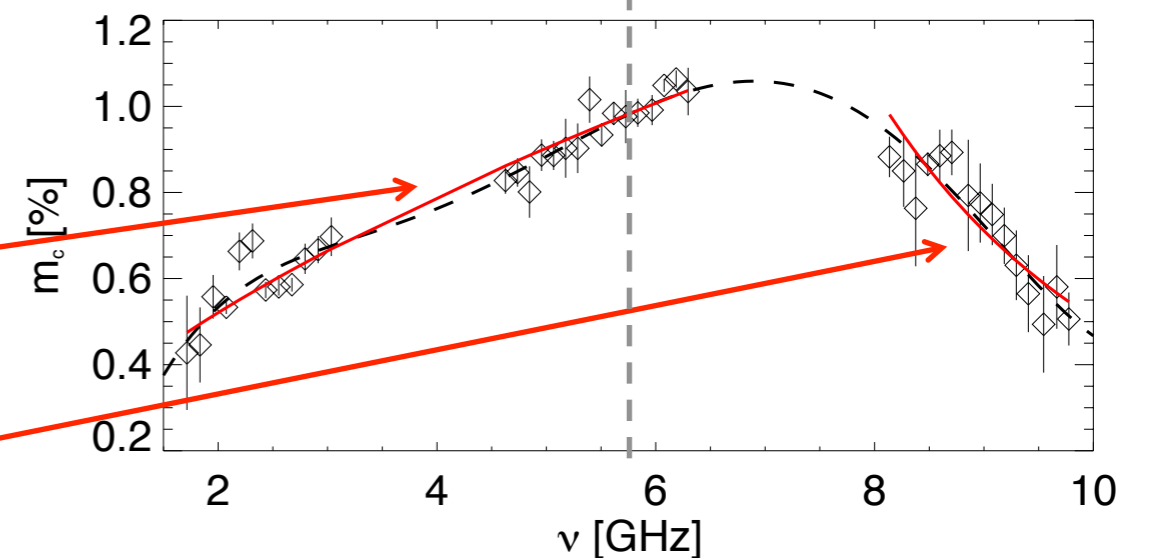
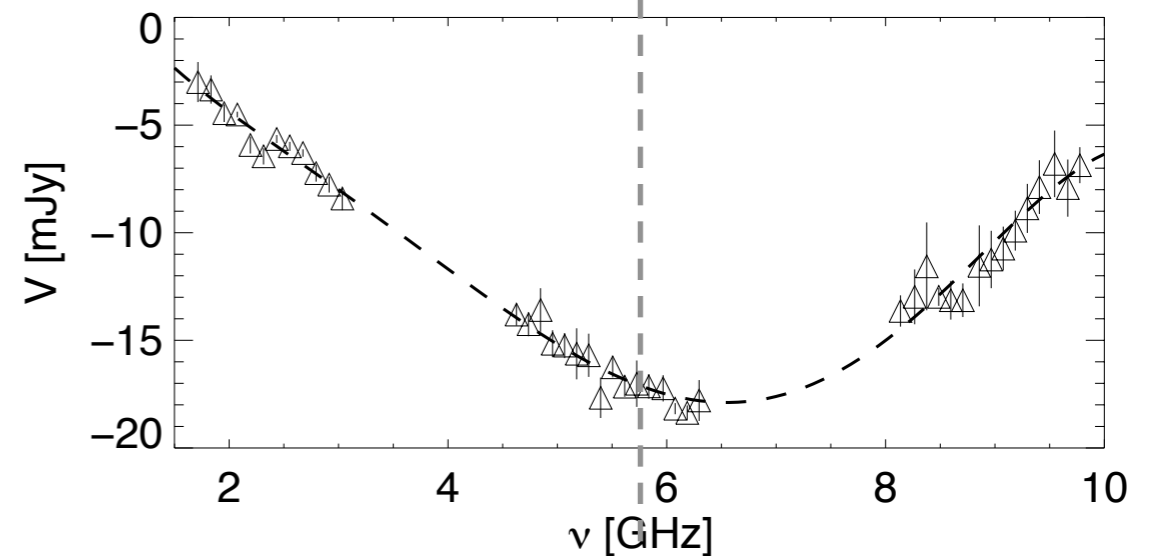
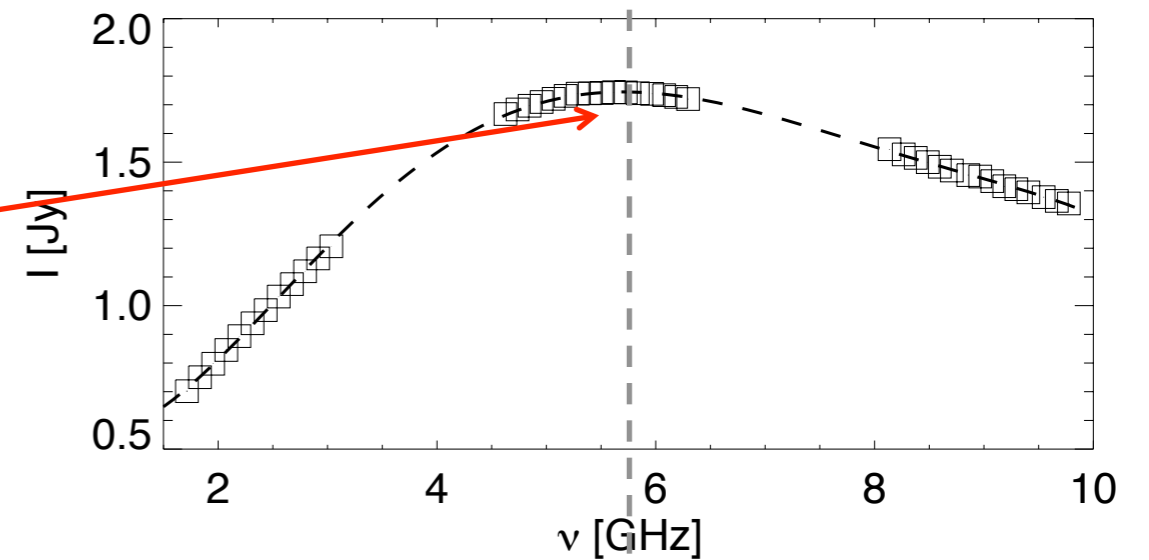
$$\nu_t \sim 5.7 \pm 0.1 \text{ GHz}$$

$$S_{\nu_t} \sim 1.744 \pm 0.001 \text{ GHz}$$

- CP increases with frequency in optically thick regime (inverted CP spectrum)
- Consistent with B&K79 conically self-similar jet
- Unity optical depth surface moves further upstream in the jet at higher frequencies, where the uniform B-field component is stronger

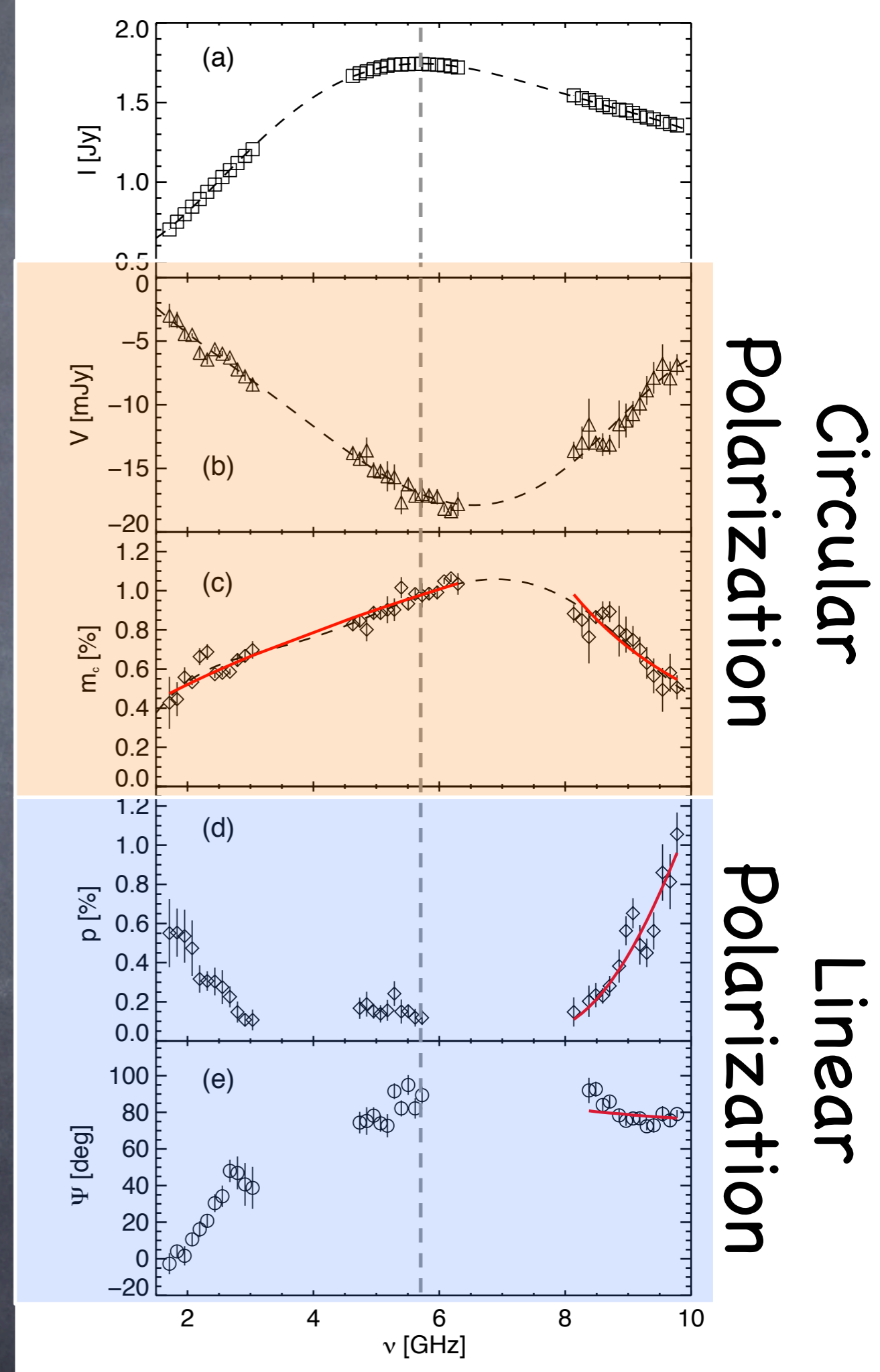
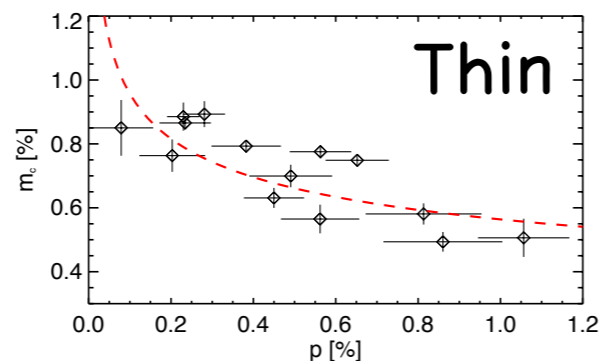
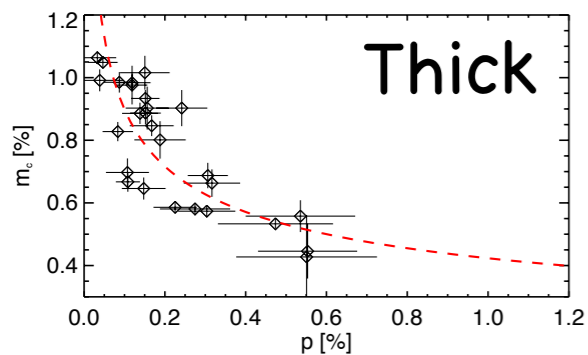
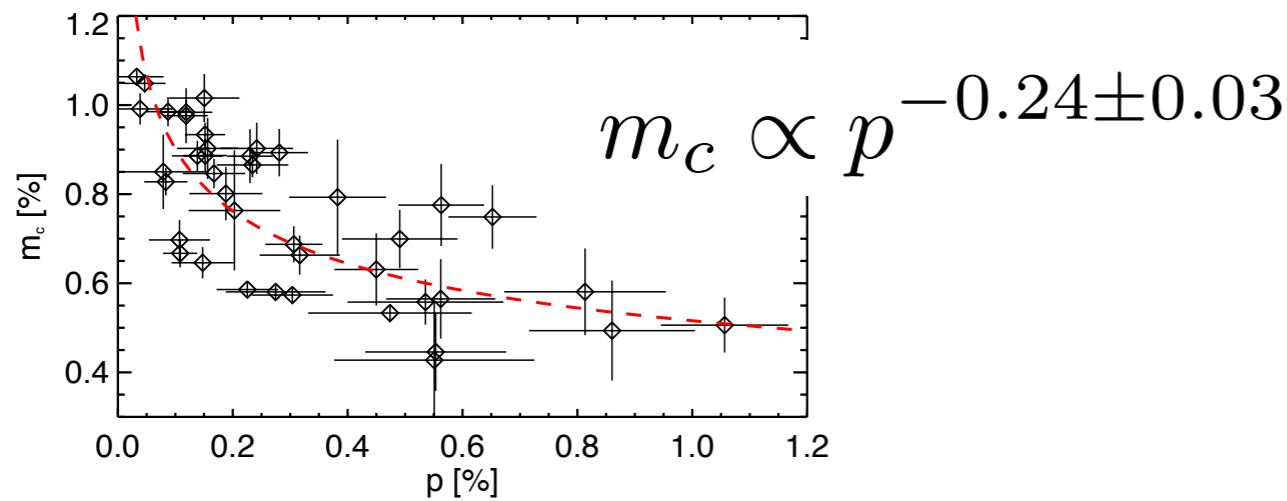
$$m_c \propto \nu^{+0.60 \pm 0.03}$$

$$m_c \propto \nu^{-3.0 \pm 0.4}$$



PKS B2126-128

- Linear and circular polarization anti-correlated
- Spearman rank test: -0.8
- Clear signature of Faraday conversion of LP to CP



PKS B2126-128

- Faraday rotation and depolarization
- Faraday thick or thin?
- Unconstrained for small range in λ^2

$$P = p e^{-2\sigma_{\text{RM}}^2 \lambda^4} e^{2i(\Psi_0 + \text{RM}\lambda^2)}$$

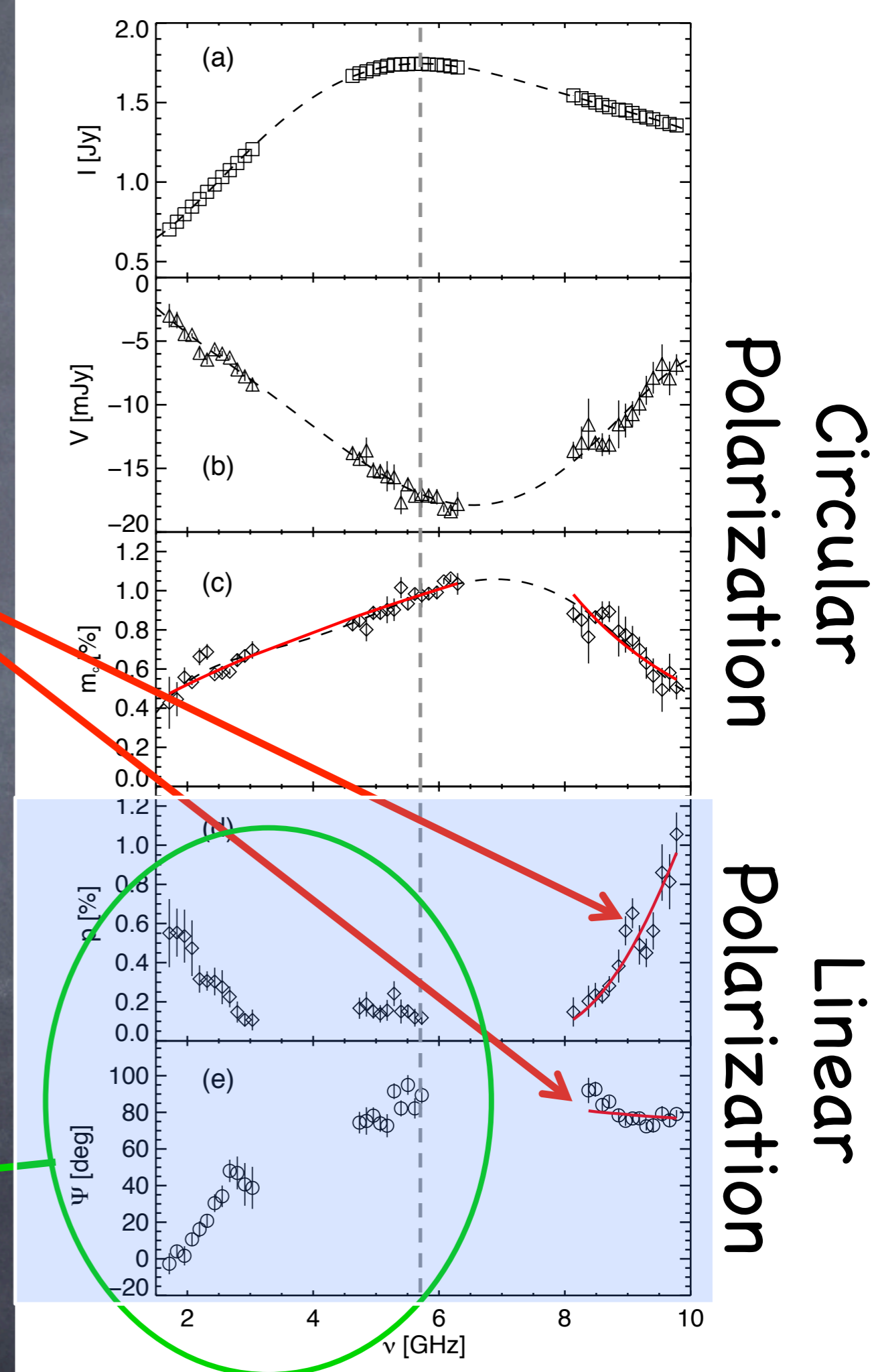
$$\text{RM} = 190 \pm 21 \text{ rad m}^{-2}$$

$$\sigma_{\text{RM}} = 1056 \pm 67 \text{ rad m}^{-2}$$

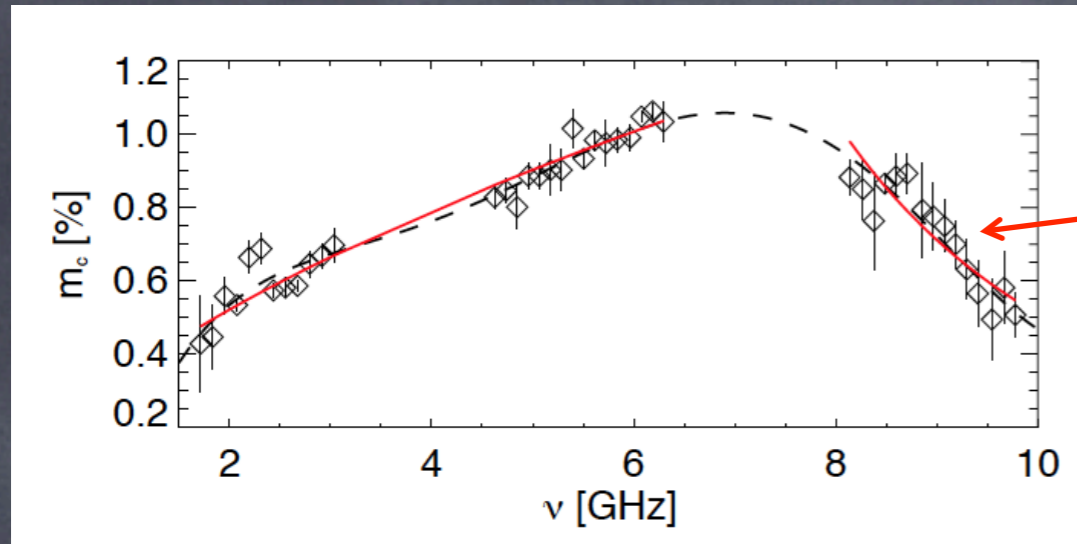
$$P = p e^{-2\sigma_{\text{RM}}^2 \lambda^4} \text{Sinc}(\text{RM}\lambda^2) e^{2i\Psi}$$

large increase in p
towards longer
wavelengths

Change in sign of RM?



RRM fit (8 – 10 GHz)



$$V \propto \sin \frac{c^3 \text{RRM}}{\nu^3}$$

$$B = (10^{-9} L^{-1} \text{RRM})^{-1/4}$$

Macquart+01, Kennett & Melrose 98

- Fitting to 8–10 GHz: $\text{RRM} = 647 \pm 12 \text{ rad/m}^3$
- $\text{RRM}_{\text{int}} = \text{RRM} ((1+z)/d)^3 = 50 \text{ rad/m}^3$ for $d=10$
- Assuming equipartition: $B = 15 \text{ mG}$ (for $L = 1 \text{ pc}$)

$$m_c^{\text{int}} \sim \left(\frac{2.8B}{\nu_{\text{obs}}(1+z)/\delta} \right)^{1/2}$$

- Intrinsic CP for completely uniform B-field: $\sim 0.3\%$ at 8 GHz

SSA analysis

$$\begin{aligned} \nu_m &= 5.7 \pm 0.1 \text{ GHz} \\ I_{\nu_m} &= 1.744 \pm 0.001 \text{ Jy} \end{aligned} \quad B = 10^{-5} b(\alpha) S_m^{-2} \nu_m^5 \theta^4 (1+z)^{-1} \delta \text{ [G]}$$

- Marscher+79,83,87
- Angular size of emission region ~ 0.5 mas (using $B = 15$ mG)

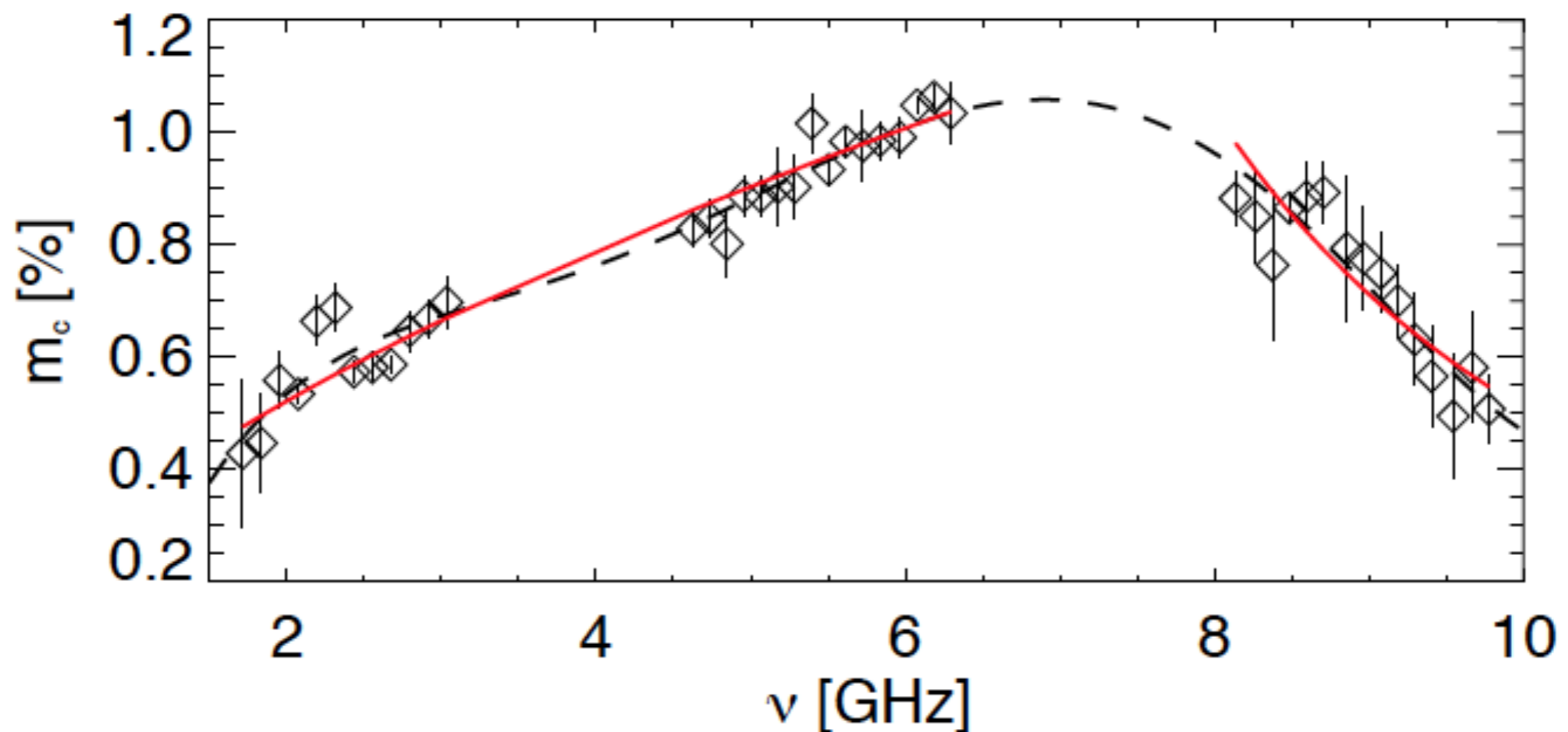
$$U_B = \frac{B^2}{8\pi} \text{ [erg cm}^{-3}\text{]}$$

$$U_{re} = 2 \ln(\nu_2/\nu_m) D^{-1} \theta^{-9} \nu_m^{-7} S_m^4 (1+z)^7 \delta^{-5} \text{ [ergs cm}^{-3}\text{]}$$

- $U_{re} \gg U_B$ (particles dominate jet energy density)
- No constraints on g_{\min} or jet composition
- VLBI observations + radiative transfer simulations required

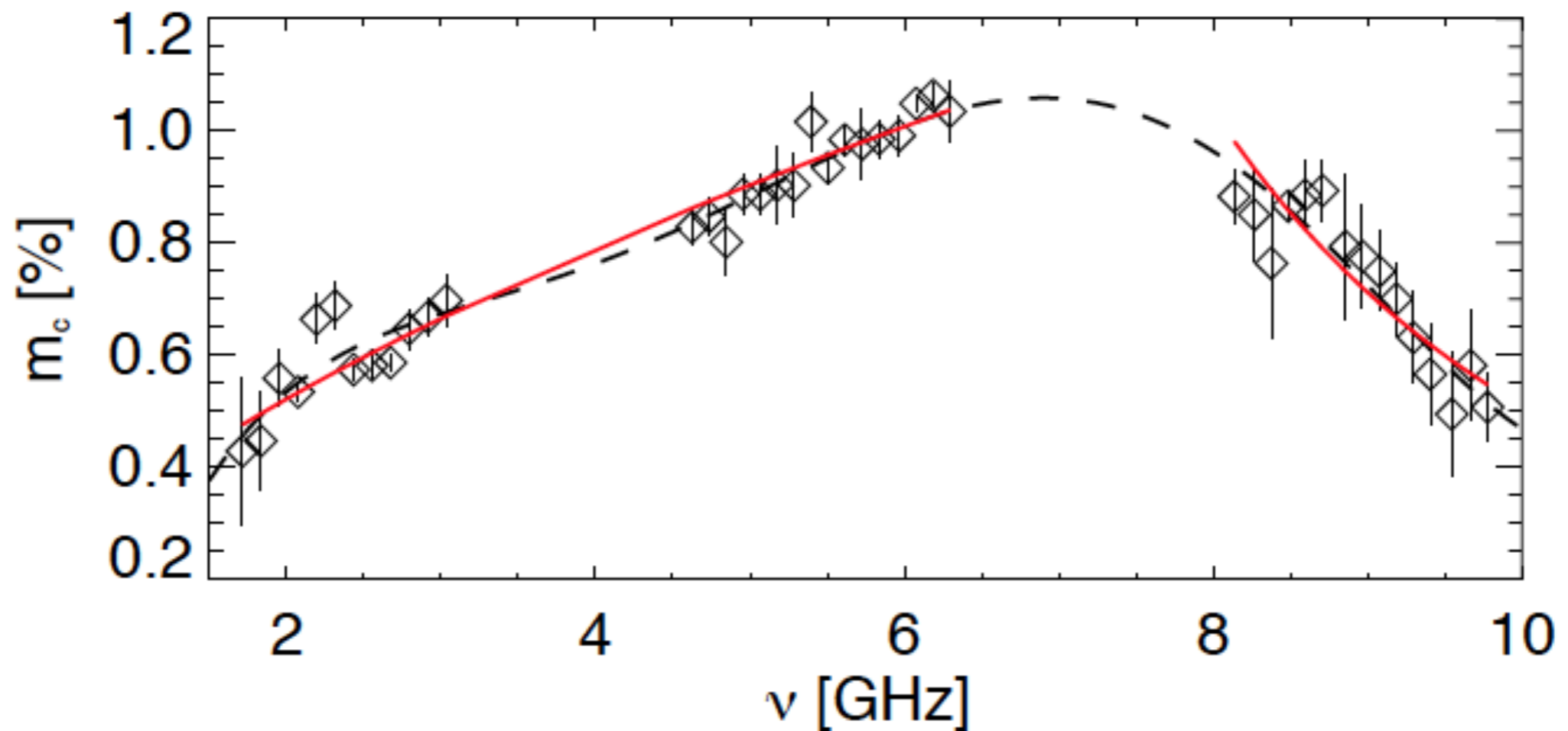
CP generation mechanisms (Homan+09)

1. Stochastic production of CP from tangled B-field: requires intrinsic or conversion or combination. CP should vary in sign and amount across different frequencies
2. Intrinsic CP from a strong (tens of mG) vector-ordered B-field (assuming all emitting particles are e-s, "e-p jet")

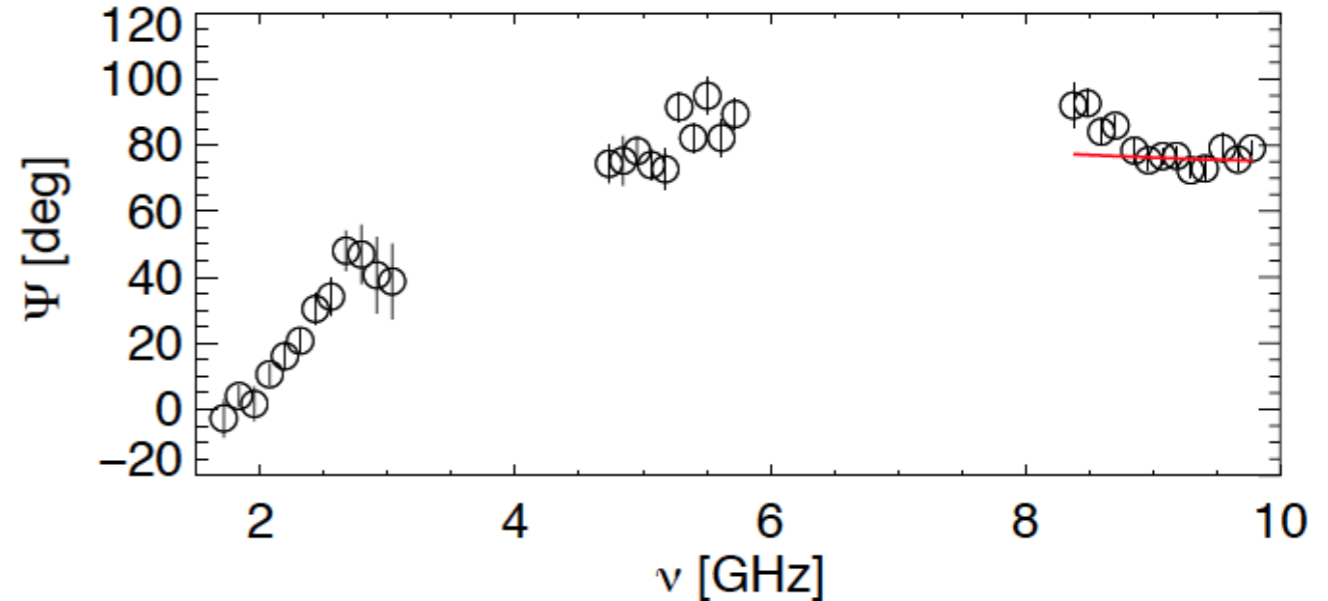
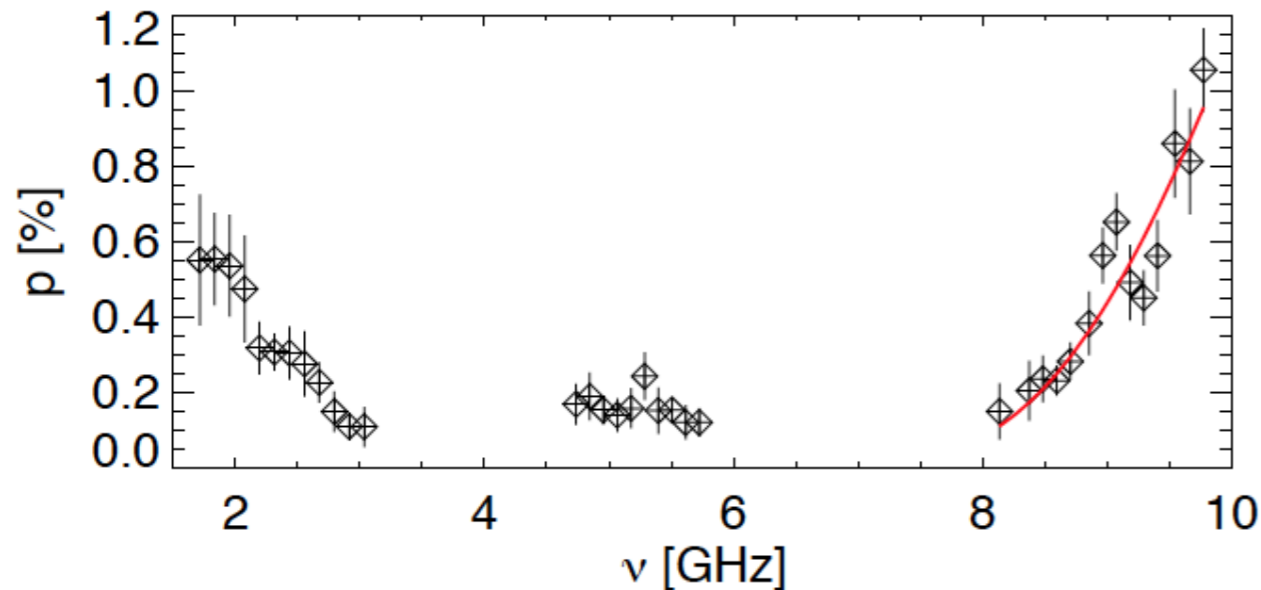


CP generation mechanisms (Homan+09)

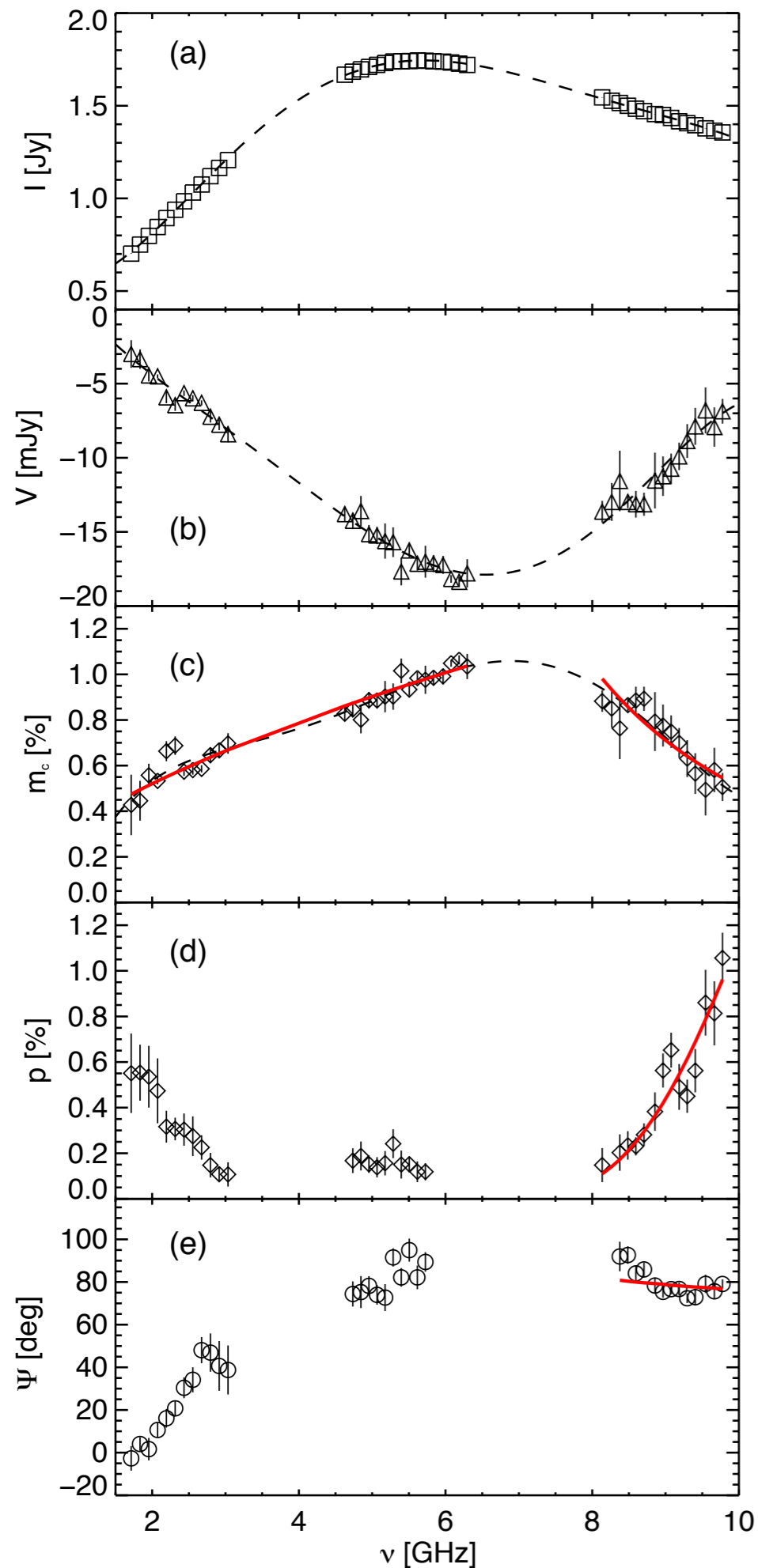
1. Stochastic generation of CP from tangled fields requires intrinsic or conversion or ν in sign and amount across different frequencies
2. Intrinsic CP (tens of μ ordered B-field (assuming all particles are ...))



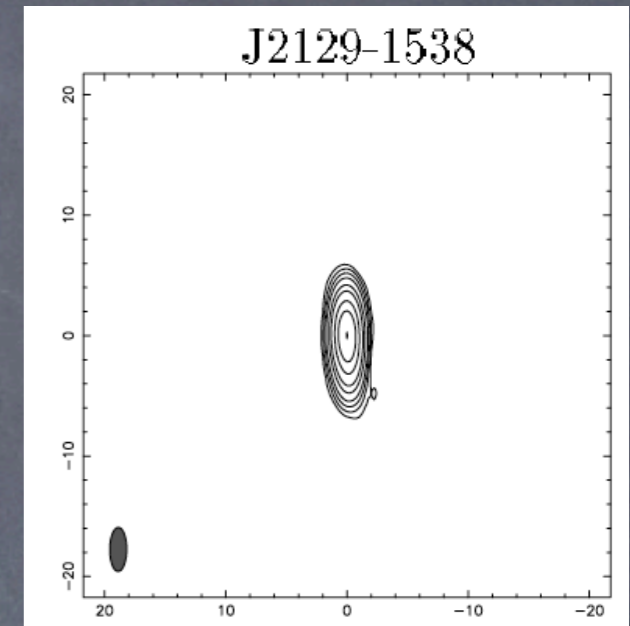
CP generation mechanisms (Homan+09)



3. Conversion of LP to CP by internal Faraday rotation (thermal or low energy e^-). **Small Faraday depth** so that B-field at front of jet converts LP from back of jet into some CP. 2 and 3 will act together to some degree.
4. **High rotation depth** version of 3, where no large scale B-field order is required and net CP can be produced in large amounts with very little net LP
5. Faraday conversion from a helical B-field. **No internal Faraday rotation** required as field orientation at back of jet is already different to that at front of jet



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- $M_{\text{BH}} \sim 10^{10} M_{\text{sol}}$
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- **Dominant CP generation mechanism: Faraday conversion of LP to CP**

- Conclusions:

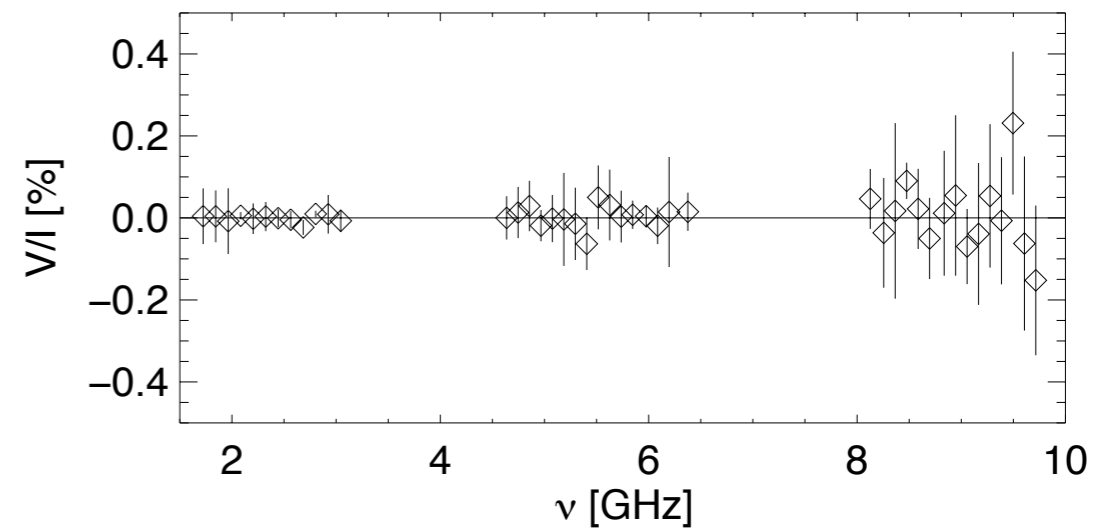
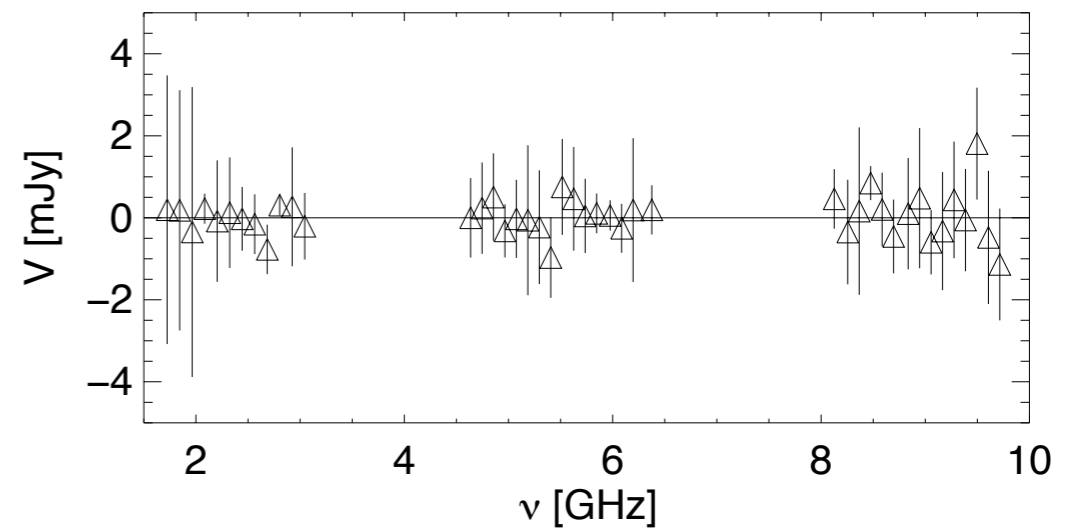
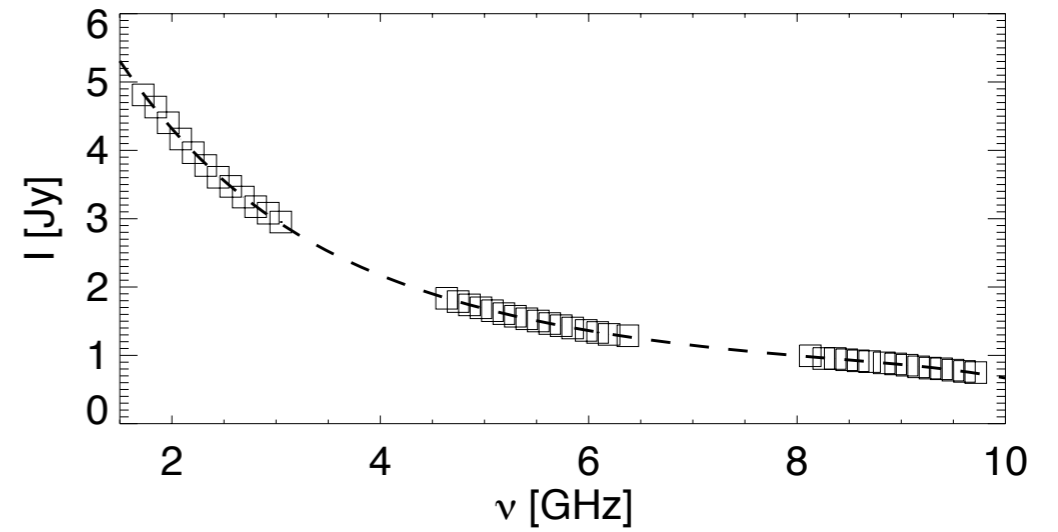
1. **little or no thermal plasma within the jet**
2. **cannot yet distinguish between magnetic twist and internal Faraday rotation models**
3. **relativistic particle content unconstrained**

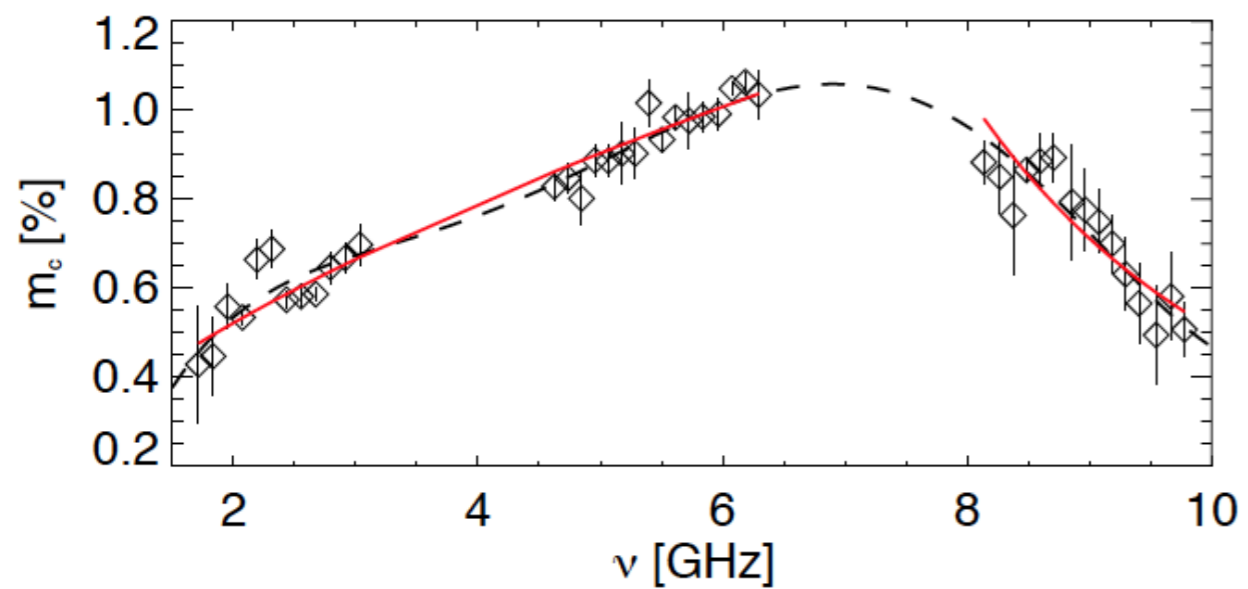
Thank You



PKS B0252-712

- Source to check calibration against
- CSS radio galaxy





Red: internal FR
Green: no internal FR

