

Poynting flux dissipation in jets

John Kirk¹

Iwona Mochol¹ Takanobu Amano²

¹Max-Planck-Institut für Kernphysik
Heidelberg, Germany

²Dept. Earth & Planetary Science, University of Tokyo

Granada, 14th June 2013

Problem

- Relativistic jets launched with high magnetization parameter: $\sigma \gg 1$.
- Collimation slow $\Rightarrow \sigma$ may stay large
- Magnetic fields are too springy — shocks do not change the Poynting flux substantially.
- Fermi I acceleration doesn't work well in magnetized $\sigma \sim 1$ shocks

Potential Solution

Dissipation is possible if magnetic field (not Lorentz factor) in jet fluctuates (e.g., twists, reversals of polarity at launch).

Potential Solution

Dissipation is possible if magnetic field (not Lorentz factor) in jet fluctuates (e.g., twists, reversals of polarity at launch).

How?

Potential Solution

Dissipation is possible if magnetic field (not Lorentz factor) in jet fluctuates (e.g., twists, reversals of polarity at launch).

How?

- 1 Wait long enough [JK & Mochol ApJ 729,104 \(2011\)](#)

Potential Solution

Dissipation is possible if magnetic field (not Lorentz factor) in jet fluctuates (e.g., twists, reversals of polarity at launch).

How?

- 1 Wait long enough [JK & Mochol ApJ 729,104 \(2011\)](#)
- 2 Hit an obstacle:

Potential Solution

Dissipation is possible if magnetic field (not Lorentz factor) in jet fluctuates (e.g., twists, reversals of polarity at launch).

How?

- 1 Wait long enough [JK & Mochol ApJ 729,104 \(2011\)](#)
- 2 Hit an obstacle:
 - In MHD description: formation of current sheets \Rightarrow compression of these at a weak shock \Rightarrow enhanced reconnection rate
Solar wind: [Drake et al, ApJ 709, 963 \(2010\)](#)
Pulsars: [Sironi & Spitkovsky, ApJ 741, 39 \(2011\)](#)

Potential Solution

Dissipation is possible if magnetic field (not Lorentz factor) in jet fluctuates (e.g., twists, reversals of polarity at launch).

How?

- 1 Wait long enough [JK & Mochol ApJ 729,104 \(2011\)](#)
- 2 Hit an obstacle:
 - In MHD description: formation of current sheets \Rightarrow compression of these at a weak shock \Rightarrow enhanced reconnection rate
Solar wind: [Drake et al, ApJ 709, 963 \(2010\)](#)
Pulsars: [Sironi & Spitkovsky, ApJ 741, 39 \(2011\)](#)
 - For *under-dense* plasmas: shock causes fluctuations to convert into electromagnetic modes forming a dissipative precursor
[Amano & Kirk ApJ 770, 18 \(2013\)](#),
[Mochol & Kirk arXiv:1303.6434](#)

Lengthscales

Why worry about under-dense plasmas?

- For an MHD description, require $\lambda \gg \lambda_g, c/\omega_p$
- In the Crab pulsar wind $r \ll 10^{-3} \times$ termination shock radius
- In a synchrotron emitting e^\pm jet

$$\lambda \gg \lambda_g = 3 \times 10^{15} \nu_{16}^{1/2} B_{\text{nT}}^{-3/2} \text{ cm}$$

Lengthscales

Why worry about under-dense plasmas?

- For an MHD description, require $\lambda \gg \lambda_g, c/\omega_p$
- In the Crab pulsar wind $r \ll 10^{-3} \times$ termination shock radius
- In a synchrotron emitting e^\pm jet

$$\lambda \gg \lambda_g = 3 \times 10^{15} \nu_{16}^{1/2} B_{\text{nT}}^{-3/2} \text{ cm}$$

Otherwise:

Lengthscales

Why worry about under-dense plasmas?

- For an MHD description, require $\lambda \gg \lambda_g, c/\omega_p$
- In the Crab pulsar wind $r \ll 10^{-3} \times$ termination shock radius
- In a synchrotron emitting e^\pm jet

$$\lambda \gg \lambda_g = 3 \times 10^{15} \nu_{16}^{1/2} B_{nT}^{-3/2} \text{ cm}$$

Otherwise:

- **Nonlinear superluminal modes** (Arka, Mochol)

Lengthscales

Why worry about under-dense plasmas?

- For an MHD description, require $\lambda \gg \lambda_g, c/\omega_p$
- In the Crab pulsar wind $r \ll 10^{-3} \times$ termination shock radius
- In a synchrotron emitting e^\pm jet

$$\lambda \gg \lambda_g = 3 \times 10^{15} \nu_{16}^{1/2} B_{nT}^{-3/2} \text{ cm}$$

Otherwise:

- **Nonlinear superluminal modes** (Arka, Mochol)
- **Electromagnetically modified shocks** (Amano)

Two-fluid simulations

Simplest description that includes electromagnetic modes is one with two charged fluids

Two-fluid simulations

Simplest description that includes electromagnetic modes is one with two charged fluids

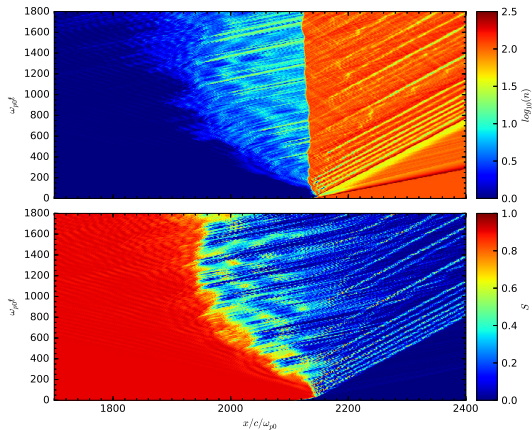
- Relativistic, finite temperature electron & positron fluids
- 1D in space, 3D in momentum and EM fields

Two-fluid simulations

Simplest description that includes electromagnetic modes is one with two charged fluids

- Relativistic, finite temperature electron & positron fluids
- 1D in space, 3D in momentum and EM fields
- Initial conditions:
 - Left half: circularly polarized, cold, static shear, $\gamma = 40$, $\sigma = 10$, $\lambda \approx \lambda_g/4$
 - Right half: shocked (R-H conditions) unmagnetized plasma

Time evolution

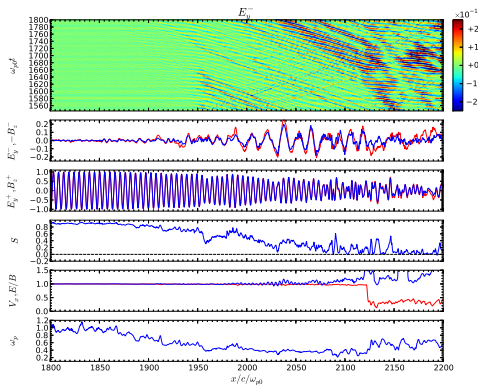


$$\Gamma = 40$$

$$\sigma = 10$$

$$\omega = 1.2\omega_p$$

Wave helicity



Positive helicity injected wave (E^+ , B^+).

Backwards propagating, negative helicity waves generated.

$E > B$ in precursor and downstream ($v_{\text{wave}} = B/E$).

Simulation Results

- Poynting flux dissipated completely
- A precursor containing strong electromagnetic waves is formed
- A hydrodynamic shock remains

Implications/Conclusions

Particle acceleration:

- In a magnetized jet, the power in fluctuations with short length scale ($\lambda < \lambda_g$) can be dissipated at an electromagnetically modified shock front
- Particle acceleration by the first order Fermi is possible at the hydrodynamic sub-shock

Implications/Conclusions

Particle acceleration:

- In a magnetized jet, the power in fluctuations with short length scale ($\lambda < \lambda_g$) can be dissipated at an electromagnetically modified shock front
- Particle acceleration by the first order Fermi is possible at the hydrodynamic sub-shock

Radiation:

- A signature from the electromagnetic precursor is possible from both *thermal* particles and accelerated particles
- Electric vector polarization angle *perpendicular* to the jet — mechanism similar to synchro-Compton (Rees 1971):
- Possibly measurable degree of circular polarization from accelerated particles penetrating the precursor